

A PROFILE DECOMPOSITION APPROACH TO THE $L_t^\infty(L_x^3)$ NAVIER-STOKES REGULARITY CRITERION

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ABSTRACT. In this paper we continue to develop an alternative viewpoint on recent studies of Navier-Stokes regularity in critical spaces, a program which was started in the recent work [14] by C. Kenig and the second author. Specifically, we prove that strong solutions which remain bounded in the space $L^3(\mathbb{R}^3)$ do not become singular in finite time, a known result established in [8] by Escauriaza, Seregin and Šverák in the context of suitable weak solutions. Here, we use the method of “critical elements” which was recently developed by Kenig and Merle to treat critical dispersive equations. Our main tool is a “profile decomposition” for the Navier-Stokes equations in critical Besov spaces which we develop here. As a byproduct of this tool, assuming a singularity-producing initial datum for Navier-Stokes exists in a critical Lebesgue or Besov space, we show there is one with minimal norm, generalizing a result of Rusin and Šverák [23].

INTRODUCTION

We consider the incompressible Navier-Stokes equations in \mathbb{R}^d ,

$$(0.1) \quad (NS) \quad \begin{cases} \frac{\partial u}{\partial t} = \Delta u - \nabla \cdot (u \otimes u) - \nabla \pi, \\ \nabla \cdot u = 0, \\ u|_{t=0} = u_0 \end{cases}$$

for $(x, t) \in \mathbb{R}^d \times (0, T)$, where $u = u(x, t)$ is the velocity vector field and $\pi(x, t)$ is the associated pressure function. For $d \geq 3$, global weak solutions are known to exist, but their uniqueness (as well as the smoothness of the solution for smooth data) has remained an open problem since the pioneering work [21]. There exist several conditional

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results, of which Serrin's criterion is perhaps the most well-known: if a weak solution u is such that

$$(0.2) \quad u \in L^p([0, T]; L^q(\mathbb{R}^d)) \text{ with } \frac{2}{p} + \frac{d}{q} = 1, \quad q > d,$$

then u is smooth on $(0, T)$. On the other hand, there is a long line of work on constructing local in time solutions, from [13] to [20]. In this framework of local in time (strong, e.g. unique) solutions, Serrin's criterion may be understood as a non blow-up criterion at time T : e.g. if u is a strong solution with $u_0 \in L^d(\mathbb{R}^d)$, that is $u \in C([0, T]; L^d(\mathbb{R}^d))$, and if (0.2) is satisfied, then one may (continuously and uniquely) extend the solution u past time T .

In the recent important work [8], Escauriaza-Seregin-Šverák obtained the endpoint version of Serrin's criterion¹: $u \in L^\infty([0, T]; L^3(\mathbb{R}^3))$ implies no blow-up; they work with the so-called suitable weak solutions introduced in [4]. A similar, but much more precise result in the (smaller) space $L^\infty([0, T]; \dot{H}^{\frac{1}{2}}(\mathbb{R}^3))$ was obtained recently by Seregin in [25]; the main point in that result is that it is proved that the $\dot{H}^{\frac{1}{2}}$ norm of the solution $u(t)$ blows up as t goes to blow-up time, and not just for a subsequence as is known in the $L^3(\mathbb{R}^3)$ case (see [24] for a partial result in that direction)².

Our goal here is to obtain such a result by a somewhat different route, following the concentration-compactness methods developed by Kenig-Merle in the context of energy critical dispersive equations ([15, 16]) and then extended to subcritical problems ([18]) or supercritical problems ([17]). In our context, the Navier-Stokes equations are supercritical with respect to their only known a priori bound, which is the energy inequality (the $L^2(\mathbb{R}^3)$ norm of a solution is decreasing), while a scale invariant norm for the data is $\dot{H}^{\frac{1}{2}}(\mathbb{R}^3)$ or $L^3(\mathbb{R}^3)$. In [14], C. Kenig and the second author carried out such a program for solutions $u \in L^\infty([0, T]; \dot{H}^{\frac{1}{2}}(\mathbb{R}^3))$. The first step in following the Kenig-Merle roadmap is to prove the existence of a so-called "critical element" (or minimal blow-up solution), which follows from suitable profile decompositions of the data, the linear solution (to the heat equation) and the nonlinear solution (to the full Navier-Stokes system). Such decompositions were introduced by P. Gérard in [11] to study the defect of compactness of the Sobolev embedding theorem, and then used by the same author and H. Bahouri to study the critical defocusing wave equation [2]. In the context of Navier-Stokes, they were developed in [10]

¹[8] treats the case $d = 3$; the case $d > 3$ was treated similarly later in [7].

²The result [24] was actually very recently improved by Seregin in [26], who obtains the blow up of $u(t)$ in the $L^3(\mathbb{R}^3)$ case.

by the first author, and served as a crucial tool in implementing the roadmap in [14]; this explains why the result in [14] applies to $\dot{H}^{\frac{1}{2}}(\mathbb{R}^3)$ rather than $L^3(\mathbb{R}^3)$, as [10] implements profile decompositions in the Sobolev scale. Recently the second author extended profile decompositions ([19]) to study the embedding $L^d(\mathbb{R}^d) \hookrightarrow \dot{B}_{p,q}^{\frac{d}{p}-1}(\mathbb{R}^d)$, where the latter are ((NS)-critical) homogeneous Besov spaces, with $p > d$, or more generally within the Besov scale itself.

Our main goal is threefold:

- we develop profile decompositions for solutions to the Navier-Stokes equations with data in $L^d(\mathbb{R}^d)$ or $\dot{B}_{p,q}^{\frac{d}{p}-1}(\mathbb{R}^d)$ with $1 \leq p, q < 2d + 3$, extending the results from [10]; the main difficulty here compared to [10] consists in handling multilinear interactions between profiles and remainders, as well as a lack of orthogonality in $L^3(\mathbb{R}^3)$ for the profile decomposition.
- we use this profile decomposition to implement the Kenig-Merle roadmap for solutions $u \in L^\infty([0, T]; L^3(\mathbb{R}^3))$. The Hilbert nature of $\dot{H}^{\frac{1}{2}}(\mathbb{R}^3)$ proved helpful in [14] to deal with weak convergence issues, as well as again with multilinear interactions. We need to face these issues here, however eventually we obtain a streamlined argument which leads to weak convergence toward zero at blow-up time for critical elements. Once this “compactness” result is proved, Serrin’s endpoint criterion is obtained as in [14], following closely the backward uniqueness argument of [8];
- we use the profile decomposition in another direction, extending recent work of Rusin-Šverák [23]: we prove that there always exists a minimal blow-up initial datum in $L^d(\mathbb{R}^d)$ if any such datum exists, and that moreover the set of such data is compact in $\dot{B}_{p,q}^{\frac{d}{p}-1}(\mathbb{R}^d)$, with $d < p \leq q \leq +\infty$, up to transformational invariance of the equations. We moreover prove a similar statement, involving two different Besov spaces in the aforementioned scale, with $p < 2d + 3$.

The next section introduces the function spaces we shall be using, and collects a few well-known facts about the Navier-Stokes system and its solutions in those function spaces. The next three sections are devoted to the profile decomposition, the regularity criterion and the minimal blow-up data, respectively.

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problem which we treat in Section 3 below.

After completion of this work, we learned of [12] where a result in the same spirit as our last section is proved, namely the existence of initial data with minimal $L^3(\mathbb{R}^3)$ norm for potential Navier-Stokes singularities; in [12] the compactness in $L^3(\mathbb{R}^3)$ up to translation-dilation is also obtained.

1. PRELIMINARIES

For the convenience of the reader, we start by recalling the usual definition of Besov spaces. We usually write X as a shorthand for the function space $X(\mathbb{R}^d)$, where d is the space dimension.

Definition 1.1. *Let ϕ be a function in \mathcal{S} such that $\widehat{\phi} = 1$ for $|\xi| \leq 1$ and $\widehat{\phi} = 0$ for $|\xi| > 2$, and define $\phi_j(x) := 2^{dj}\phi(2^jx)$. Then the frequency localization operators are defined by*

$$S_j := \phi_j * \cdot, \quad \Delta_j := S_{j+1} - S_j.$$

Let f be in \mathcal{S}' . We say f belongs to $\dot{B}_{p,q}^s$ if

- The partial sum $\sum_{-m}^m \Delta_j f$ converges to f as a tempered distribution if $s < \frac{d}{p}$ and after taking the quotient with polynomials if not.
- The sequence $\epsilon_j := 2^{js} \|\Delta_j f\|_{L^p}$ belongs to ℓ^q , and its ℓ^q -norm defines the Besov norm of f .

We shall also need a slight modification of those spaces, taking into account the time variable.

Definition 1.2. *Let $u(x, t) \in \mathcal{S}'(\mathbb{R}^{d+1})$ and let Δ_j be a frequency localization with respect to the x variable. We shall say that u belongs to $\mathcal{L}^\rho([a, b]; \dot{B}_{p,q}^s)$ if $u(t) \in \dot{B}_{p,q}^s$ for all $t \in [a, b]$ and*

$$2^{js} \|\Delta_j u\|_{L^\rho([a,b]; L_x^p)} =: \epsilon_j \in \ell^q.$$

We define

$$\|u\|_{\mathcal{L}^\rho([a,b]; \dot{B}_{p,q}^s)} := \|2^{js} \|\Delta_j u\|_{L^\rho([a,b]; L_x^p)}\|_{\ell^q},$$

and $\mathcal{L}_T^\rho(\dot{B}_{p,q}^s) := \mathcal{L}^\rho([0, T]; \dot{B}_{p,q}^s)$.

Remark 1.3. *In the case where $\rho \geq q$ one has of course the embedding $\mathcal{L}^\rho([a, b]; \dot{B}_{p,q}^s) \hookrightarrow L^\rho([a, b]; \dot{B}_{p,q}^s)$ due to Minkowski's inequality.*

Let us introduce some notation. For any p in $[1, \infty)$ we define

$$s_p := -1 + \frac{d}{p}.$$

For any initial datum $u_0 \in \dot{B}_{p,q}^{s_p}$, with $d < p \leq q < +\infty$, we shall denote by $NS(u_0)$ the local in time strong solution to the Navier-Stokes equation (0.1). For clarity, by “solution” to (0.1) in the strong (sometimes called “mild”) sense, we mean a divergence-free solution u to

$$u_t = \Delta u - \mathbb{P}\nabla \cdot (u \otimes u), \quad u|_{t=0} = u_0$$

(equivalent to solving (0.1) for the “right” π in the Duhamel sense, where \mathbb{P} is the projection operator onto divergence-free vector fields. Such solutions were obtained in [5] for $3 < p \leq 6$ and for all $p < +\infty$ in [22], and we refer to the appendix of [9] for a proof which is tailored to our purposes. The specific case of $L^d(\hookrightarrow \dot{B}_{p,q}^{-(1-\frac{d}{p})})$ data is included in such a result, as any additional “regularity” is propagated along the flow (see again [9] for a proof of this well-known fact).

Let us make those results on the Cauchy problem more precise. We define the function space

$$E_{p,q}(T) := \mathcal{L}^\infty([0, T]; \dot{B}_{p,q}^{s_p}) \cap \mathcal{L}^{\frac{2p}{p+1}}([0, T]; \dot{B}_{p,q}^{s_p+1+\frac{1}{p}}).$$

In particular, $E_{p,q} \subset \mathcal{L}^r(\dot{B}_{p,q}^{s_p+\frac{2}{r}})$ for $\frac{2p}{p+1} \leq r \leq \infty$ by interpolation. We recall (see e.g. [9]) that $NS(u_0)$ belongs to $E_{p,q}(T)$ for some time T , and one may define a maximal time $T^* = T^*(u_0)$ such that this holds for any $T < T^*$ (actually the solution belongs to $\mathcal{L}^\infty([0, T]; \dot{B}_{p,q}^{s_p}) \cap \mathcal{L}^1([0, T]; \dot{B}_{p,q}^{s_p+2})$ but that fact will not be used here). If the initial datum is small enough then $T^* = \infty$ (and under such a condition one may include $q = \infty$, although one cannot in general obtain local solutions for $q = \infty$). Moreover, u belongs to $E_{p,q}(T^*)$ if and only if $T^* = \infty$, and in that case one has (see [9]) that $\lim_{t \rightarrow \infty} \|NS(u_0)(t)\|_{\dot{B}_{p,q}^{s_p}} = 0$ as well. Finally recall that if $NS(u_0)$ belongs to $E_{p,q}(T)$ and if u_0 belongs to $\dot{B}_{a,b}^{s_a}$ (resp. $L^d(\mathbb{R}^d)$) with $a \leq p$ and $b \leq q$, then $NS(u_0)$ belongs to $E_{a,b}(T)$ (resp. $C([0, T]; L^d(\mathbb{R}^d))$) with the same life span (see [9] for instance, or [1]).

2. PROFILE DECOMPOSITIONS

2.1. Notation and statement of the result. In what follows, we shall need the following notion, where the dimension d is always chosen such that $d \geq 2$.

Definition 2.1. For $j \in \mathbb{N}$, let $\{(\lambda_{j,n}, x_{j,n})\}_{n=1}^\infty \subset (0, \infty) \times \mathbb{R}^d$ be a sequence of “scales” $\lambda_{j,n}$ and “cores” $x_{j,n}$. We say that such a set of

sequences is (pairwise) orthogonal if

$$(2.1) \quad j \neq j' \implies \left\{ \begin{array}{l} \text{either} \quad \lim_{n \rightarrow +\infty} \frac{\lambda_{j,n}}{\lambda_{j',n}} + \frac{\lambda_{j',n}}{\lambda_{j,n}} = +\infty \\ \\ \text{or} \\ \\ \frac{\lambda_{j,n}}{\lambda_{j',n}} \equiv 1 \quad \text{and} \quad \lim_{n \rightarrow +\infty} \frac{|x_{j,n} - x_{j',n}|}{\lambda_{j,n}} = +\infty . \end{array} \right.$$

Let us define, for functions $f \in L^d(\mathbb{R}^d)$, norm-invariant transformations depending on translation and scaling parameters $x_{j,n}$ and $\lambda_{j,n}$ respectively by

$$(2.2) \quad \Lambda_{j,n}f(x) := \frac{1}{\lambda_{j,n}} f\left(\frac{x - x_{j,n}}{\lambda_{j,n}}\right) .$$

Then the reason for the terminology in the previous definition becomes clear by considering the following quantity:

$$(2.3) \quad \int_{\mathbb{R}^d} |\Lambda_{j,n}f(x)|^{d-1} |\Lambda_{j',n}g(x)| \, dx \\ = \frac{\lambda_{j,n}}{\lambda_{j',n}} \int_{\mathbb{R}^d} |f(y)|^{d-1} \left| g\left(\frac{\lambda_{j,n}}{\lambda_{j',n}}y + \frac{x_{j,n} - x_{j',n}}{\lambda_{j',n}}\right) \right| \, dy .$$

One similarly has equality if on the right-hand side one interchanges j and j' and the arguments of the functions, and hence such a term tends to zero as $n \rightarrow \infty$ if (2.1) holds, since we may approximate f and g in $L^d(\mathbb{R}^d)$ by \mathcal{C}_0^∞ -functions. This implies that

$$\|\Lambda_{j,n}f + \Lambda_{j',n}g\|_{L^d(\mathbb{R}^d)}^d = \|\Lambda_{j,n}f\|_{L^d(\mathbb{R}^d)}^d + \|\Lambda_{j',n}g\|_{L^d(\mathbb{R}^d)}^d + o(1)$$

as $n \rightarrow \infty$ if the scales are orthogonal. Therefore $\Lambda_{j,n}f$ and $\Lambda_{j',n}g$ are ‘‘asymptotically orthogonal’’ in $L^d(\mathbb{R}^d)$ (and similar statements will be shown to hold in other critical spaces as well).

We first recall two theorems which were proved³ in [19]. The first one deals with the defect of compactness of the embedding $L^d \hookrightarrow \dot{B}_{p,q}^{s_p}$ with $d < p, q \leq +\infty$.

Theorem 1 ([19]). *Let $p, q \in \mathbb{R}$ satisfy $d < p, q \leq +\infty$. Let $\{\varphi_n\}_{n=1}^\infty$ be a bounded sequence in $L^d(\mathbb{R}^d)$ and let ϕ_0 be any weak limit point of $\{\varphi_n\}_n$. Then, after possibly replacing $\{\varphi_n\}_n$ by a subsequence which*

³It is known that the wavelet-basis characterization of scalar function-spaces used in [19] extend as well to vector-fields (and in fact one may use divergence-free wavelets, see e.g. [3]). Moreover, the slightly more specific formulations we give here are a simple consequence of the theorems in [19].

we relabel $\{\varphi_n\}_n$, there exists a sequence of profiles $\{\phi_j\}_{j=1}^\infty$ of $L^d(\mathbb{R}^d)$, and for each $j \in \mathbb{N}$ sequences $\{(\lambda_{j,n}, x_{j,n})\}_{n=1}^\infty$ in $(0, \infty) \times \mathbb{R}^d$ of scales and cores which are orthogonal in the sense of Definition 2.1 such that, for all $n, J \in \mathbb{N}$, if we define ψ_n^J by

$$(2.4) \quad \varphi_n(x) = \phi_0(x) + \sum_{j=1}^J \frac{1}{\lambda_{j,n}} \phi_j \left(\frac{x - x_{j,n}}{\lambda_{j,n}} \right) + \psi_n^J(x)$$

the following properties hold:

- the function ψ_n^J is a remainder in the sense that

$$(2.5) \quad \lim_{J \rightarrow \infty} \left(\limsup_{n \rightarrow \infty} \|\psi_n^J\|_{\dot{B}_{p,q}^{s_p}} \right) = 0;$$

- there is a norm⁴ $\|\cdot\|_{L^d(\mathbb{R}^d)}$ which is equivalent to $\|\cdot\|_{L^d(\mathbb{R}^d)}$ such that

$$(2.6) \quad \sum_{j=1}^\infty \|\phi_j\|_{L^d(\mathbb{R}^d)}^d \leq \liminf_{n \rightarrow \infty} \|\varphi_n\|_{L^d(\mathbb{R}^d)}^d$$

and, for each integer J ,

$$(2.7) \quad \|\psi_n^J\|_{L^d(\mathbb{R}^d)} \leq \|\varphi_n\|_{L^d(\mathbb{R}^d)} + o(1) \quad \text{as } n \rightarrow \infty;$$

- for any integer j , the following properties hold: either $\lambda_{j,n} = 1$ and $\lim_{n \rightarrow \infty} |x_{j,n}| = +\infty$, or $\lim_{n \rightarrow \infty} \lambda_{j,n} \in \{0, +\infty\}$.

The second statement deals with the defect of compactness of the embedding $\dot{B}_{a,b}^{s_a} \hookrightarrow \dot{B}_{p,q}^{s_p}$ with $1 \leq a < p \leq +\infty$ and $1 \leq b \leq (p/a)b \leq q \leq +\infty$.

Theorem 2 ([19]). *Let $a, b, p, q \in [1, \infty]$ satisfy $1 \leq a < p \leq +\infty$ and $1 \leq b \leq (p/a)b \leq q \leq +\infty$. Let $\{\varphi_n\}_{n=1}^\infty$ be a bounded sequence in $\dot{B}_{a,b}^{s_a}(\mathbb{R}^d)$, and let ϕ_0 be any weak limit point of $\{\varphi_n\}$. Then, after possibly replacing $\{\varphi_n\}_n$ by a subsequence which we relabel $\{\varphi_n\}_n$, there exists a sequence of profiles $\{\phi_j\}_{j=1}^\infty$ in $\dot{B}_{a,b}^{s_a}$, and for each integer $j \geq 1$ a sequence $\{(\lambda_{j,n}, x_{j,n})\}_{n=1}^\infty$ of scales and cores which are orthogonal in the sense of Definition 2.1 such that, for all $n, J \in \mathbb{N}$, if we define ψ_n^J by*

$$(2.8) \quad \varphi_n(x) = \phi_0(x) + \sum_{j=1}^J \frac{1}{\lambda_{j,n}} \phi_j \left(\frac{x - x_{j,n}}{\lambda_{j,n}} \right) + \psi_n^J(x)$$

the following properties hold:

⁴This norm is based on a wavelet basis expansion for functions in Besov and Triebel-Lizorkin spaces.

- the function ψ_n^J is a remainder in the sense that

$$(2.9) \quad \lim_{J \rightarrow \infty} \left(\limsup_{n \rightarrow \infty} \|\psi_n^J\|_{\dot{B}_{p,q}^{s_p}} \right) = 0;$$

- there is a norm $\|\cdot\|_{\dot{B}_{a,b}^{s_a}}$ which is equivalent to $\|\cdot\|_{\dot{B}_{a,b}^{s_a}}$ such that for each integer $n \in \mathbb{N}$, denoting $\tau := \max\{a, b\}$, one has

$$(2.10) \quad \left\| \left(\|\phi_j\|_{\dot{B}_{a,b}^{s_a}} \right)_{j=0}^{\infty} \right\|_{\ell^\tau} \leq \liminf_{n' \rightarrow \infty} \|\varphi_{n'}\|_{\dot{B}_{a,b}^{s_a}}$$

and, for each integer J ,

$$(2.11) \quad \|\psi_n^J\|_{\dot{B}_{a,b}^{s_a}} \leq \|\varphi_n\|_{\dot{B}_{a,b}^{s_a}} + o(1) \quad \text{as } n \rightarrow \infty;$$

- for any integer j , the following properties hold: either $\lambda_{j,n} = 1$ and $\lim_{n \rightarrow \infty} |x_{j,n}| = +\infty$, or $\lim_{n \rightarrow \infty} \lambda_{j,n} \in \{0, +\infty\}$.

It should be noted (see e.g. the proof of (3.2) below with $s = 0$) that the ϕ_j 's are weak limit points of the sequence

$$\lambda_{j,n} \varphi_n (\lambda_{j,n} \cdot + x_{j,n}).$$

Consequently if Theorem 1 or Theorem 2 is applied to a sequence of divergence free vector fields, each profile produced by the theorem is also a divergence free vector field.

We shall now apply these results to a sequence of bounded initial data to (0.1). The following statement is the analogue of Theorem 2 of [10] in the context of critical Besov spaces. For the sake of simplicity, we shall only consider the case $p = q$; the restriction on p is technical in nature but it allows to deal with $L^3(\mathbb{R}^3)$ data, and we plan to address this issue in future work.

Theorem 3 (NSE Evolution of Profile Decompositions). *Suppose $d \geq 3$ and fix $a, b, p \in \mathbb{R}$ satisfying $d < a \leq b < p < 2d + 3$. Let $\{\varphi_n\}_{n=1}^{\infty}$ be a bounded sequence of divergence-free vector fields in $\dot{B}_{a,b}^{s_a}(\mathbb{R}^d)$, and let ϕ_0 be any weak limit point of $\{\varphi_n\}$. Then, after possibly relabeling the sequence due to the extraction of a subsequence following an application of Theorem 2 (of which we retain the same notation for the profile decomposition), defining $u_n := NS(\varphi_n)$ and $U_j := NS(\phi_j) \in E_{a,b}(T_j)$ for every integer $j \in \mathbb{N}$ (where T_j is any real number smaller than the life span T_j^* of ϕ_j , and $T_j = \infty$ if $T_j^* = \infty$), the following properties hold:*

- there is a finite (possibly empty) subset I of \mathbb{N} such that⁵

$$\forall j \in I, \quad T_j < \infty \quad \text{and} \quad \forall j \in \mathbb{N} \setminus I, \quad U_j \in E_{a,b}(\infty).$$

⁵This is equivalent to $T_j^* < \infty \iff j \in I$ by [9] and the small-data theory.

Moreover setting $\tau_n := \min_{j \in I} \lambda_{j,n}^2 T_j$ if I is nonempty and $\tau_n := \infty$ otherwise, we have

$$\sup_n \|u_n\|_{E_{a,b}(\tau_n)} < \infty.$$

- if there exists a time $T \in \mathbb{R}^+ \cup \{+\infty\}$ such that $\{u_n\}$ is uniformly bounded in $E_{a,b}(T)$, then if I is nonempty we must have

$$(2.12) \quad \forall n \in \mathbb{N}, \quad T < \min_{j \in I} \lambda_{j,n}^2 T_j^*,$$

and therefore in such a case the scales of concentration for which $\lim_{n \rightarrow \infty} \lambda_{j,n} = 0$ (small scales) generate global solutions of (NS) (i.e., the corresponding $T_j^* = +\infty$).

- finally there exists some large $J_0 \in \mathbb{N}$ such that for each $J > J_0$, there exists $N(J) \in \mathbb{N}$ such that for all $n > N(J)$, all $t \leq \tau_n$ and all $x \in \mathbb{R}^d$, setting $w_n^J := e^{t\Delta}(\psi_n^J)$ and defining r_n^J by

$$(2.13) \quad u_n(x, t) = U_0(x, t) + \sum_{j=1}^J \frac{1}{\lambda_{j,n}} U_j \left(\frac{x - x_{j,n}}{\lambda_{j,n}}, \frac{t}{\lambda_{j,n}^2} \right) + w_n^J(x, t) + r_n^J(x, t),$$

then w_n^J and r_n^J are small remainders in the sense that

$$(2.14) \quad \lim_{J \rightarrow \infty} \left(\limsup_{n \rightarrow \infty} \|w_n^J\|_{E_{p,p}(\infty)} \right) = \lim_{J \rightarrow \infty} \left(\limsup_{n \rightarrow \infty} \|r_n^J\|_{E_{p,p}(\tau_n)} \right) = 0 ;$$

Remark 2.2. As a consequence of the proof of Theorem 3, the decomposition (2.13) actually holds for all indices J (once the theorem is proved as stated, the remainder r_n^J for $1 \leq J \leq J_0$ may be defined by (2.13) which now holds up to time τ_n) but it should be clear that such a decomposition is of interest mostly when enough profiles have been extracted, yielding a small remainder.

Remark 2.3. Although the first step in proving Theorem 3 is to apply Theorem 2, if moreover $\{\varphi_n\}_{n=1}^\infty$ is a bounded sequence of divergence-free vector fields in $L^d(\mathbb{R}^d)$, one may instead first apply Theorem 1 and standard embeddings to see that the sequence enjoys all the properties in the conclusion of Theorem 1 as well as those in the conclusion of Theorem 3.

Remark 2.4. The bound $p < 2d + 3$ is directly related to Lemma 2.5 below; we did not seek the optimal values (w.r.t. the original a, b), and for $a = b = d$, the regularity $s_p + 2/p$ which appears in this Lemma may be taken positive by choosing $p < d + 2$. This in turn would avoid the direct use of paraproduct estimates in the proof of Lemma 2.7, but

we feel worth pointing out that negative regularity can be allowed in Lemma 2.5. At any rate, paraproduct techniques are at the heart of the estimates in [9] or any of the product estimates, both of which we heavily rely on.

2.2. Proof of Theorem 3. The first step consists of course in applying Theorem 2 (or Theorem 1, if applicable) to the sequence of initial data. We briefly comment on the choice of indices: our main goal being to deal with a sequence of data in $L^3(\mathbb{R}^3)$, a natural choice (in spatial dimension d) is $a = b = d$, for which in the profile decomposition of the data we may pick any $p = q$ close to d (hence even less than $d + 2$, see Remark 2.4). For general a, b , we may relabel $a = b = \sup(a, b)$ (using Sobolev's embedding), and then use Theorem 2 for $p = q$ strictly larger than a . Thus we can always use $p = q = \sup(a, b) + \varepsilon$ as our indices for the remainder space. Now let us continue with the proof.

With the notation of Theorem 2 we define

$$u_n := NS(\varphi_n), \quad U_j := NS(\phi_j) \in E_{a,b}(T_j) \quad \text{and} \quad w_n^J := e^{t\Delta}(\psi_n^J).$$

Then due to (2.9) and standard linear heat estimates we have

$$(2.15) \quad \lim_{J \rightarrow \infty} \left(\limsup_{n \rightarrow \infty} \|w_n^J\|_{E_{p,p}(\infty)} \right) = 0.$$

Moreover due to the orthogonality property (2.10), the sequence $\{\phi_j\}$ goes to zero in the space $\dot{B}_{a,b}^{s_a}$ as j goes to infinity. This implies that there is j_0 such that for all $j > j_0$, there is a global unique solution associated with ϕ_j , as $\|\phi_j\|_{\dot{B}_{a,b}^{s_a}} < \varepsilon_0$ (the smallness constant of small data theory). Hence, I will be a subset of $\{0, \dots, j_0\}$ which proves the first part of the first statement in Theorem 3.

All other statements will follow from obtaining careful bounds on both profiles and r_n^J , by taking advantage of the local Cauchy theory and its perturbed variants (as set up e.g. in [9]; see the appendix for a slightly more general statement).

By the local Cauchy theory we can solve the Navier-Stokes equation (0.1) with data φ_n for each integer n , and produce a unique solution $u_n \in E_{a,b}(T_{u_n})$ for some time $T_{u_n} < T^*(\varphi_n)$. Now let us define, for any $J \geq 0$

$$r_n^J(x, t) := u_n(x, t) - \left(\sum_{j=0}^J \Lambda_{j,n} U_j(x, t) + w_n^J(x, t) \right),$$

where $\Lambda_{0,n}U_0(x,t) := U_0(x,t)$, and we abuse our earlier notation for dilations/translations to define

$$\Lambda_{j,n}U_j(x,t) := \frac{1}{\lambda_{j,n}}U_j\left(\frac{x-x_{j,n}}{\lambda_{j,n}}, \frac{t}{\lambda_{j,n}^2}\right).$$

To be consistent, from now on we define $\lambda_{0,n} := 1$ and $x_{0,n} := 0$, and remark that the lifetime of the re-scaled U_j has become at least $\lambda_{j,n}^2 T_j$. Therefore, the function $r_n^J(x, \cdot)$ is defined a priori for $t \in [0, t_n]$, where

$$t_n = \min\left(T_{u_n}; \min_{j \in I} \lambda_{j,n}^2 T_j; \infty\right) = \min(\tau_n; T_{u_n})$$

with the notation of Theorem 3. Our main goal consists in proving that r_n^J is actually defined on $[0, \tau_n]$ (at least if J is large enough), which will be a consequence of the perturbation theory for the Navier-Stokes equation. In the process, we shall obtain the desired uniform limiting property

$$\lim_{J \rightarrow \infty} \left(\limsup_{n \rightarrow \infty} \|r_n^J\|_{E_{p,p}(\tau_n)} \right) = 0.$$

Let us write the equation satisfied by r_n^J . It turns out to be easier to write that equation after a re-scaling in space-time. For convenience, let us re-order the functions $\Lambda_{j,n}U_j$, for $0 \leq j \leq j_0$, in such a way that, for some $n_0 = n_0(j_0)$ sufficiently large, we have

$$(2.16) \quad \forall n \geq n_0, \quad j \leq j' \leq j_0 \implies \lambda_{j,n}^2 T_j^* \leq \lambda_{j',n}^2 T_{j'}^*$$

(some of these terms may equal infinity), where T_j^* is the maximal life span of ϕ_j (such a reordering is possible on a fixed and finite number of profiles due to the orthogonality of scales). In particular, with this ordering we have $\tau_n = \lambda_{0,n}^2 T_0$, and we note that $\lambda_{j,n}^2 T_j^*$ is the life span of $\Lambda_{j,n}U_j$.

The inverse of our dilation/translation operator $\Lambda_{j,n}$ is

$$(2.17) \quad \Lambda_{j,n}^{-1}f(s,y) := \lambda_{j,n}f(\lambda_{j,n}y + x_{j,n}, \lambda_{j,n}^2 s).$$

We then define, for every integer J ,

$$\forall j \leq J, \quad U_n^{j,0} := \Lambda_{0,n}^{-1}\Lambda_{j,n}U_j, \quad R_n^{J,0} := \Lambda_{0,n}^{-1}r_n^J,$$

$$W_n^{J,0} := \Lambda_{0,n}^{-1}w_n^J \quad \text{and} \quad U_n^0 := \Lambda_{0,n}^{-1}u_n.$$

Clearly we have

$$R_n^{J,0}(s,y) := U_n^0(s,y) - \left(\sum_{j=0}^J U_n^{j,0}(s,y) + W_n^{J,0}(s,y) \right),$$

and $R_n^{J,0}$ (which for the time being is defined for times s in $[0, t_n^0)$ where $t_n^0 := \min\{T_0, \lambda_{0,n}^{-2} T_{u_n}\}$) solves the following system:

$$(2.18) \quad \begin{cases} \partial_s R_n^{J,0} + \mathbb{P}(R_n^{J,0} \cdot \nabla R_n^{J,0}) - \Delta R_n^{J,0} + Q(R_n^{J,0}, F_n^{J,0}) & = G_n^{J,0} \\ R_n^{J,0}|_{s=0} & = 0, \end{cases}$$

where we recall that $\mathbb{P} := Id - \nabla \Delta^{-1}(\nabla \cdot)$ is the projection onto divergence free vector fields,

$$Q(a, b) := \mathbb{P}(a \cdot \nabla b + b \cdot \nabla a)$$

for two vector fields a, b , and, finally, where

$$(2.19) \quad F_n^{J,0} := \sum_{j \leq J} U_n^{j,0} + W_n^{J,0},$$

and

$$(2.20) \quad G_n^{J,0} := -\frac{1}{2} \sum_{\substack{j \neq j' \\ (j, j') \in \{0, \dots, J\}^2}} Q(U_n^{j,0}, U_n^{j',0}) \\ - \sum_{j \leq J} Q(U_n^{j,0}, W_n^{J,0}) - \frac{1}{2} Q(W_n^{J,0}, W_n^{J,0}).$$

Notice here that in re-ordering the profiles, we may have lost the fact that $\lambda_{0,n} = 1$; however we have a (simpler notation for a) uniform lower bound on the lifetime of all possibly blowing-up profiles: for any real number T_0 smaller than T_0^* , (2.16) gives

$$\forall j \in \mathbb{N}, \quad \{U_n^{j,0}\}_{n \geq n_0} \text{ is bounded in } E_{a,b}(T_0).$$

However, we have no uniform control over the sum $F_n^{J,0}$ which enters the drift term in the perturbed equation (2.18). In order to use perturbative bounds on this system, as stated in Proposition A.1 in the appendix, we need such a good control on the drift term (which will come from orthogonality arguments), and smallness of the forcing term $G_n^{J,0}$.

We start with the drift term. Notice that we do not claim uniform boundedness in $E_{p,p}$ but rather in a weaker space, which will be enough for our purposes (the reader might notice that we could replace p in the following statement by $b (< p)$ but this will not be necessary and would make notations even more cumbersome).

Lemma 2.5. *The sequence $(F_n^{J,0})$ is bounded in $\mathcal{L}^p([0, T_0]; \dot{B}_{p,p}^{s_p + \frac{2}{p}})$, uniformly in J .*

Let us prove this lemma. Notice that

$$F_n^{J,0} = \sum_{j \leq j_0} U_n^{j,0} + \sum_{j_0+1}^J U_n^{j,0} + W_n^{J,0}$$

and by the small data theory, $U_n^{j,0}$ (for $j > j_0$) and $W_n^{J,0}$ are bounded in $E_{p,p}(\mathbb{R}^+)$ by their respective initial data in $\dot{B}_{p,p}^{s_p}$; in particular, for all $1 \leq r \leq +\infty$,

$$(2.21) \quad \forall j > j_0, \quad \|U_n^{j,0}\|_{\mathcal{L}^r(\mathbb{R}^+; \dot{B}_{p,p}^{s_p + \frac{2}{r}})} = \|U_j\|_{\mathcal{L}^r(\mathbb{R}^+; \dot{B}_{p,p}^{s_p + \frac{2}{r}})} \lesssim \|\phi_j\|_{\dot{B}_{p,p}^{s_p}}.$$

Therefore, (2.10) yields

$$\sum_{j=j_0+1}^J \|U_n^{j,0}\|_{\mathcal{L}_t^p \dot{B}_{p,p}^{s_p + \frac{2}{p}}}^p \leq \sum_{j=j_0+1}^J \|U_n^{j,0}\|_{E_{p,p}(+\infty)}^p < +\infty,$$

where the bound is uniform in J .

We then need to use the orthogonality of cores/scales in the form of the following lemma:

Lemma 2.6. *Fix any $r \in [p, \infty)$. There exists some $\epsilon(J, n) : \mathbb{N} \times \mathbb{N} \rightarrow \mathbb{R}^+$ satisfying, for each fixed J ,*

$$\lim_{n \rightarrow \infty} \epsilon(J, n) = 0$$

and such that, for any $J' < J$,

$$\left\| \sum_{j=J'}^J U_n^{j,0} \right\|_{\mathcal{L}_t^r \dot{B}_{p,p}^{s_p + \frac{2}{r}}}^p \leq \sum_{j=J'}^J \|U_n^{j,0}\|_{\mathcal{L}_t^r \dot{B}_{p,p}^{s_p + \frac{2}{r}}}^p + \epsilon(J, n).$$

Postponing the proof of this for a moment, let us use it to complete the proof of Lemma 2.5. Lemma 2.6, along with the small data theory for Navier-Stokes, implies that

$$(2.22) \quad 2^{-p} \left\| \sum_{j=j_0+1}^J U_n^{j,0} + W_n^{J,0} \right\|_{\mathcal{L}_t^p \dot{B}_{p,p}^{s_p + \frac{2}{p}}}^p \lesssim \sum_{j=j_0+1}^{\infty} \|\phi_j\|_{B_{p,p}^{s_p}}^p + \epsilon(J, n) + \|\psi_n^J\|_{B_{p,p}^{s_p}}^p.$$

The first $j_0 + 1$ terms are then dealt with because the scaling we chose is such that the norm of $U_n^{j,0}$ in $E_{p,p}([0, T_0])$ is bounded uniformly in n , by a constant depending on T_0 , and that concludes the proof of the bound on $(F_n^{J,0})$: Lemma 2.5 is proved. \square

Now let us prove Lemma 2.6. Recall that for any $\sigma \in \mathbb{R}$, we have the following equivalence of norms, where we replace the 2^j scale by a

continuous parameter $\tau \sim 2^{-2j}$ (which is easier to handle with rescaling) and the frequency localization Δ_j by the derivative of heat kernel $K(\tau) := \tau \partial_\tau e^{\tau \Delta}$:

$$(2.23) \quad \|f\|_{\dot{B}_{p,q}^\sigma} \sim \left\| \left\| \tau^{-\sigma/2} K(\tau) f \right\|_{L^p} \right\|_{L^q(\mathbb{R}^+, \frac{d\tau}{\tau})}.$$

There is no difficulty in adding the time norm, and hence, setting

$$\gamma := -1 - ps_p/2 - p/r$$

with $1 \leq r \leq +\infty$, we have

$$(2.24) \quad \|f\|_{\mathcal{L}_t^r(\dot{B}_{p,p}^{s_p+\frac{2}{r}})}^p \sim \int_0^\infty \tau^\gamma \|K(\tau) f\|_{L_t^r L_x^p}^p d\tau.$$

We proceed now with the lemma. Writing

$$\left\| K(\tau) \sum_\ell f_\ell \right\|_{L_t^r L_x^p}^p = \left\| \left| \sum_\ell K(\tau) f_\ell \right|^p \right\|_{L_t^{\frac{r}{p}} L_x^1},$$

and applying the elementary inequality

$$(2.25) \quad \left| \left| \sum_{\ell=1}^L A_\ell \right|^p - \sum_{\ell=1}^L |A_\ell|^p \right| \leq C(L) \sum_{\ell \neq \ell'} |A_\ell| |A_{\ell'}|^{p-1},$$

to the sum inside the norm on the right which we take to be $\sum_{j=j_0+1}^J K(\tau) U_n^{j,0}$

and applying the triangle inequality in $L_t^{\frac{r}{p}} L_x^1$ (this is where we need the restriction $r \geq p$) we get $C(J)$ cross-terms like

$$\begin{aligned} & \int_0^\infty \tau^\gamma \left\| K(\tau) \left[\frac{1}{\lambda_{1,n}} U_1 \left(\frac{\cdot - x_{1,n}}{\lambda_{1,n}}, t \right) \right] (x) \right. \\ & \quad \times \left. \left(K(\tau) \left[\frac{1}{\lambda_{2,n}} U_2 \left(\frac{\cdot - x_{2,n}}{\lambda_{2,n}}, t \right) \right] (x) \right)^{p-1} \right\|_{L_t^{\frac{r}{p}} L_x^1}^p d\tau = \\ & \quad \int_0^\infty \tau^\gamma \left\| \frac{1}{\lambda_{1,n}} [K(\lambda_{1,n}^{-2} \tau) [U_1(t)]] \left(\frac{x - x_{1,n}}{\lambda_{1,n}} \right) \right. \\ & \quad \times \left. \left(\frac{1}{\lambda_{2,n}} [K(\lambda_{2,n}^{-2} \tau) [U_2(t)]] \left(\frac{x - x_{2,n}}{\lambda_{2,n}} \right) \right)^{p-1} \right\|_{L_t^{\frac{r}{p}} L_x^1}^p d\tau. \end{aligned}$$

Using the change of variables $y = \frac{x - x_{1,n}}{\lambda_{1,n}}$, $s = \lambda_{1,n}^{-2}\tau$, we see that this term equals

$$\left(\frac{\lambda_{1,n}}{\lambda_{2,n}}\right)^{\frac{2+d}{p'}} \int_0^\infty \left\| V_1(y, s) \left(V_2 \left(\frac{\lambda_{1,n}}{\lambda_{2,n}} y + \frac{x_{1,n} - x_{2,n}}{\lambda_{2,n}}, \left(\frac{\lambda_{1,n}}{\lambda_{2,n}}\right)^2 s \right) \right)^{p-1} \right\|_{L_t^{\frac{r}{p}} L_y^1} ds,$$

where $\frac{1}{p} + \frac{1}{p'} = 1$ and $V_i(y, s) = s^{\frac{\gamma}{p}} [K(s)[U_i(t)]](y)$ for $i \in \{1, 2\}$. Since $U_i(t) \in E_{p,p}$, we can approximate V_1 and V_2 in $L_s^p L_t^r L_x^p$ by smooth functions of (s, t, y) with compact support in $(0, +\infty)^2 \times \mathbb{R}^d$. It is therefore clear (by dislocation of the inner L_y^1) that the above term tends to zero as $n \rightarrow \infty$ if $\frac{\lambda_{1,n}}{\lambda_{2,n}} \rightarrow 0$ as $n \rightarrow \infty$, or if $\lambda_{1,n} \equiv \lambda_{2,n}$ and $\left|\frac{x_{1,n} - x_{2,n}}{\lambda_{2,n}}\right| \rightarrow \infty$ as $n \rightarrow \infty$. Similarly, if we appropriately chose the new variables in terms of $\lambda_{2,n}$ and $x_{2,n}$ instead, we can easily show that the term also tends to zero if $\frac{\lambda_{2,n}}{\lambda_{1,n}} \rightarrow 0$, and the result follows in view of (2.1). Lemma 2.6 is proved. \square

We now turn to the source term and prove another lemma.

Lemma 2.7. *Let $G_n^{J,0}$ be the source term defined by (2.20) and*

$$F_{p,p}(T_0) := \mathcal{L}^{\frac{2p}{p+1}}([0, T_0]; \dot{B}_{p,p}^{s_p-1+\frac{1}{p}}) + \mathcal{L}^{p'}([0, T_0]; \dot{B}_{p,p}^{s_p-\frac{2}{p}}),$$

where $1/p + 1/p' = 1$. Then

$$\lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|G_n^{J,0}\|_{F_{p,p}(T_0)} = 0.$$

First by standard product laws in Besov spaces (joint with a Hölder estimate in time) we gather that

$$\begin{aligned} \|Q(W_n^{J,0}, W_n^{J,0})\|_{\mathcal{L}^{p'}([0, T_0]; \dot{B}_{p,p}^{s_p-\frac{2}{p}})} &\leq \|W_n^{J,0} \otimes W_n^{J,0}\|_{\mathcal{L}^{p'}([0, T_0]; \dot{B}_{p,p}^{s_p+1-\frac{2}{p}})} \\ &\lesssim \|W_n^{J,0}\|_{\mathcal{L}^{2p'}([0, T_0]; \dot{B}_{p,p}^{s_p+\frac{1}{p'}})}^2. \end{aligned}$$

Note that the condition $0 < s_p + 1/p' = (d-1)/p < d/p$ allows to justify the product law.

Then by definition of $W_n^{J,0}$ and due to the scaling invariance of the space $\mathcal{L}^{2p'}([0, T_0]; \dot{B}_{p,p}^{s_p+\frac{1}{p'}})$ we obviously have

$$(2.26) \quad \lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|Q(W_n^{J,0}, W_n^{J,0})\|_{\mathcal{L}^{p'}([0, T_0]; \dot{B}_{p,p}^{s_p-2/p})} = 0.$$

Define $fg = \mathcal{T}_f g + \mathcal{T}_g f + \Pi(f, g)$ the paraproduct decomposition of the product fg , and $\zeta(f, g) := fg - \mathcal{T}_f g$. We notice that (abusing notations

as the Q entries are vectors)

$$\|Q(f, g)\|_{\mathcal{L}^{p'}([0, T_0]; \dot{B}_{p,p}^{-2+s_p+\frac{2}{p'}})} \lesssim \|fg\|_{\mathcal{L}^{p'}([0, T_0]; \dot{B}_{p,p}^{-1+s_p+\frac{2}{p'}})},$$

and we have

$$\begin{aligned} \|\zeta(\sum_{j < J} U_n^{j,0}, W_n^{J,0})\|_{\mathcal{L}^{p'}([0, T_0]; \dot{B}_{p,p}^{-1+s_p+\frac{2}{p'}})} &\lesssim \|W_n^{J,0}\|_{\mathcal{L}^r([0, T_0]; \dot{B}_{p,p}^{s_p+\frac{2}{r}})} \\ &\times \left\| \sum_{j < J} U_n^{j,0} \right\|_{\mathcal{L}^p([0, T_0]; \dot{B}_{p,p}^{s_p+\frac{2}{p}})} \end{aligned}$$

where $1/p' = 1/r + 1/p$. The product law is allowed because the sum of the two respective regularities is again $2s_p + 2/r + 2/p = 2(d-1)/p > 0$. Notice that the regularity of $W_n^{J,0}$ is $\sigma = d/p + (p-4)/p > d/p$ for large p , which explains why we subtracted the paraproduct term carrying the corresponding low frequencies of f . This specific term is handled differently, as

$$\begin{aligned} \|\mathcal{T}(\sum_{j < J} U_n^{j,0}) W_n^{J,0}\|_{\mathcal{L}^{\frac{2p}{p+1}}([0, T_0]; \dot{B}_{p,p}^{s_p-1+\frac{p+1}{p}})} &\lesssim \|W_n^{J,0}\|_{\mathcal{L}^{\frac{2p}{p-1}}([0, T_0]; \dot{B}_{p,p}^{s_p+\frac{p-1}{p}})} \\ &\times \left\| \sum_{j < J} U_n^{j,0} \right\|_{\mathcal{L}^p([0, T_0]; \dot{B}_{p,p}^{s_p+\frac{2}{p}})} \end{aligned}$$

where we do not care whether the sum of regularities $-1 + (2d+1)/p$ is negative, thanks to the frequency distribution in the paraproduct.

The profiles $U_n^{j,0}$ in $\mathcal{L}^p([0, T_0]; \dot{B}_{p,p}^{s_p+\frac{2}{p}})$ are obviously under control for $j > j_0$, as seen in the proof of Lemma 2.5 (see (2.21)), and so is their sum due to Lemma 2.6. It follows that

$$(2.27) \quad \lim_{J \rightarrow \infty} \limsup_{n \rightarrow \infty} \|Q(\sum_{j < J} U_n^{j,0}, W_n^{J,0})\|_{E_{p,p}(T_0)} = 0.$$

Now we are left with the terms $Q(U_n^{j,0}, U_n^{j',0})$ for $j \neq j'$. Again, we resort to the orthogonality of the profiles to conclude: here we use the fact that, at fixed $j < J$, $U_n^{j,0}$ belongs to $E_{p,p}(T_0)$, but with no uniform bound with respect to J ; by scaling we have (with r to be specified later)

$$\|Q(U_n^{j,0}, U_n^{j',0})\|_{\mathcal{L}_{T_0}^r(\dot{B}_{p,p}^{-2+s_p+\frac{2}{r}})} = \|Q(\Lambda_{j,n} U_j, \Lambda_{j',n} U_{j'})\|_{\mathcal{L}_{(\lambda_n^0)^2 T_0}^r(\dot{B}_{p,p}^{-2+s_p+\frac{2}{r}})}$$

and we recall that

$$\Lambda_{j,n} U_j(x, t) := \frac{1}{\lambda_{j,n}} U_j\left(\frac{x - x_{j,n}}{\lambda_{j,n}}, \frac{t}{\lambda_{j,n}^2}\right)$$

is defined at least on $[0, \lambda_{0,n}^2 T_0]$ by construction.

Suppose to simplify that $T_j \leq T_{j'}$ (if not exchange j and j'). By density (for $r < \infty$) for any $\varepsilon > 0$ one can find smooth compactly supported functions F^j and $F^{j'}$ such that

$$\|F^j - U_j\|_{\mathcal{L}^{2r}([0, T_j]; \dot{B}_{p,p}^{s_p + \frac{1}{r}})} + \|F^{j'} - U^{j'}\|_{\mathcal{L}^{2r}([0, T_j]; \dot{B}_{p,p}^{s_p + \frac{1}{r}})} \leq \varepsilon.$$

Then we decompose

$$Q(U_n^{j,0}, U_n^{j',0}) = Q(\rho_n^j, \rho_n^{j'}) + Q(f_n^j, \rho_n^{j'}) + Q(\rho_n^j, f_n^{j'}) + Q(f_n^j, f_n^{j'})$$

where

$$f_n^j := \frac{1}{\lambda_{j,n}} F^j \left(\frac{x - x_{j,n}}{\lambda_{j,n}}, \frac{t}{\lambda_{j,n}^2} \right)$$

and $U_n^{j,0} = f_n^j + \rho_n^j$. The same argument as above enables us to write that for all j, j' ,

$$(2.28) \quad \sup_n \|Q(\rho_n^j, \rho_n^{j'})\|_{\mathcal{L}^r([0, T_0]; \dot{B}_{p,p}^{-2+s_p+\frac{2}{r}})} \leq C\varepsilon^2$$

and

$$(2.29) \quad \sup_n \|Q(f_n^j, \rho_n^{j'}) + Q(\rho_n^j, f_n^{j'})\|_{\mathcal{L}^r([0, T_0]; \dot{B}_{p,p}^{-2+s_p+\frac{2}{r}})} \leq C\varepsilon,$$

where C only depends on T_0 (neither on j nor on j'). So finally we are left with $Q(f_n^j, f_n^{j'})$, or more accurately with the product $f_n^j f_n^{j'}$. But that product is dealt with exactly as the cross terms in the proof of Lemma 2.6, as smoothness lets us estimate the product in the space $\mathcal{L}^r([0, T_0]; \dot{B}_{r,r}^{s_r + \frac{2}{r} - 1})$ or as in Proposition 2.1 of [10], up to a suitable modification of the time exponents to reach $1 < r$.

Chosing $r = p'$, by Sobolev's embedding this implies in particular that

$$(2.30) \quad \lim_{n \rightarrow \infty} \|Q(f_n^j, f_n^{j'})\|_{\mathcal{L}^{p'}([0, T_0]; \dot{B}_{p,p}^{s_p - \frac{2}{p}})} = 0.$$

Putting together estimates (2.26) to (2.30) ends the proof of Lemma 2.7. \square

Using Lemma 2.5, we have that the drift term $F_n^{J,0}$ is uniformly bounded in $\mathcal{L}^p([0, T_0]; \dot{B}_{p,p}^{s_p + \frac{2}{p}})$ by a constant C_0 (depending on the profiles, on j_0 and on T_0), provided that at fixed J , n is chosen larger than some $N(J)$. Then if one picks J large enough, the forcing term is small enough (with respect to the smallness condition in the perturbation result recalled in Proposition A.1), and we obtain the desired bound on $R_n^{J,0}$ in $E_{p,p}(T_0)$ for $J > J_0$ and $n > N(J_0)$ thanks to Proposition A.1. This bound immediately implies that u_n is bounded in $E_{p,p}(T_0)$, and by persistence of regularity (see e.g. [9]), we get $u_n \in E_{a,b}(T_0)$.

All other statements in Theorem 3 follow easily, as in [10], which ends the proof. \square

3. SERRIN'S ENDPPOINT REGULARITY CRITERION

3.1. Preliminaries and statement of the main result. Let $X = X(\mathbb{R}^d)$ be a critical space with respect to the Navier-Stokes scaling, that is

$$\|\lambda f(\lambda \cdot)\|_X = \|f\|_X$$

for any $f \in X$, and such that local in time strong solutions can be constructed. Let $T^*(u_0)$ be the maximal time of existence of the solution $u = NS(u_0)$ in X for any $u_0 \in X$. The type of statement we would like to address is the following:

Statement 3.1. *For any $u_0 \in X$,*

$$\sup_{t \in [0, T^*(u_0))} \|NS(u_0)(t)\|_X < \infty \quad \implies \quad T^*(u_0) = +\infty .$$

Such a statement was proved in [8] for $X = L^3(\mathbb{R}^3)$ (in the context of Leray-Hopf weak solutions). An alternative proof of that theorem was given in [14] for $X = \dot{H}^{\frac{1}{2}}(\mathbb{R}^3)$ using the profile decompositions in [10]. (Of course that result is implied as well by [8].) Our goal now is to give a proof similar to the one in [14] using the profile decompositions in this article for the case $X = L^3(\mathbb{R}^3)$. This will give a different proof of the following, which was proved in [8] and also extended to $d > 3$ in [7]:

Theorem 4 (Endpoint regularity criterion). *For any $u_0 \in L^3(\mathbb{R}^3)$,*

$$\sup_{t \in [0, T^*(u_0))} \|NS(u_0)(t)\|_{L^3(\mathbb{R}^3)} < \infty \quad \implies \quad T^*(u_0) = +\infty .$$

Note that due to the time-continuity in $L^3(\mathbb{R}^3)$ of strong solutions, the left-hand side is equivalent to $NS(u_0) \in L^\infty((0, T^*(u_0)); L^3(\mathbb{R}^3))$, or in the notation of [8], $NS(u_0) \in L_{3,\infty}(0, T^*(u_0))$.

3.2. Proof of Theorem 4. Theorem 4 follows from Theorems 5 to 7 below, which will be proved in the subsequent sections. In the following, we define $A_c > 0$ by

$$A_c := \sup \left\{ A > 0 ; \sup_{t \in [0, T^*(u_0))} \|NS(u_0)(t)\|_{L^3(\mathbb{R}^3)} \leq A \implies T^*(u_0) = \infty \right\} .$$

Note that A_c is well-defined by small-data results. Moreover, if A_c is finite, then we have

$$A_c = \inf \left\{ \sup_{t \in [0, T^*(u_0))} \|NS(u_0)(t)\|_{L^3(\mathbb{R}^3)} ; T^*(u_0) < \infty \right\} .$$

Theorem 5 (Existence of a critical element). *Suppose $A_c < +\infty$, and let $u_{0,n}$ be any bounded sequence in L^3 such that $T^*(u_{0,n}) < +\infty$ and such that $A_c \leq A_n$, where*

$$A_n := \sup_{t \in [0, T^*(u_{0,n}))} \|NS(u_{0,n})(t)\|_{L^3} < \infty ,$$

and satisfying $A_n \rightarrow A_c$ as $n \rightarrow \infty$. Let $U_j = NS(\phi_j)$ be the Navier-Stokes profiles associated to $\{u_{0,n}\}$. Then there exists an integer $j_0 \in \mathbb{N}$ such that

$$T^*(\phi_{j_0}) < +\infty \quad \text{and} \quad \sup_{t \in [0, T^*(\phi_{j_0}))} \|U_{j_0}(t)\|_{L^3} = A_c .$$

We shall call any solution with the properties of U_{j_0} above a “critical element” — that is, those solutions whose supremum in L^3 is A_c and who blow up in finite time. For the next two statements we fix u_0 in L^3 . The first gives a kind of compactness property for critical elements:

Theorem 6 (Compactness of critical elements). *Suppose that A_c is finite and that $T^*(u_0) < +\infty$ and assume*

$$\sup_{t \in [0, T^*(u_0))} \|u(t)\|_{L^3} = A_c$$

where $u = NS(u_0)$. Then $u(t) \rightarrow 0$ in \mathcal{S}' as $t \nearrow T^*(u_0)$.

Finally, we state the following, which corresponds to the “rigidity theorem” in [14]:

Theorem 7 (Global existence of critical elements). *Assume that the solution $u := NS(u_0)$ satisfies*

$$\sup_{t \in [0, T^*(u_0))} \|u(t)\|_{L^3} < +\infty$$

and moreover $u(t) \rightarrow 0$ in \mathcal{S}' as $t \nearrow T^*(u_0)$. Then $T^*(u_0) = +\infty$.

Theorems 5 to 7 immediately imply that $A_c = +\infty$, and Theorem 4 is proved. Indeed, if $A_c < +\infty$ then Theorem 5 guarantees us some j_0 such that $u_0 = \phi_{j_0}$ satisfies the conditions of Theorem 6, but then Theorem 7 contradicts the fact that $T^*(u_0) < +\infty$. \square

In order to prove Theorems 5 to 7, we shall need the following result, which was proved in [9]:

Theorem 8 ([9]). *Let u_0 be some divergence-free vector field in X such that $T^*(u_0) = +\infty$, where X is either $L^3(\mathbb{R}^3)$ or $\dot{B}_{p,p}^{s_p}(\mathbb{R}^3)$. Then*

$$\lim_{t \rightarrow +\infty} \|NS(u_0)(t)\|_X = 0 .$$

We shall also need the following lemma in the spirit of Lemma 2.6 whose proof we outline below (see the proof of (3.18) in [14] for more details):

Lemma 3.2. *Suppose $\sup_n \|u_{0,n}\|_{L^3(\mathbb{R}^3)} < \infty$ and let $\{U_j\}$ be the associated Navier-Stokes profiles given by Theorem 3. Let $\{t_n\} \subset \mathbb{R}^+$ be any sequence such that $t_n \leq \tau_n$ for all n . There exists a subsequence in n such that the following is true: for any $\epsilon > 0$ and $J \in \mathbb{N}$, there exists $N_0 = N_0(J, \epsilon) \in \mathbb{N}$ such that, for any $J' < J$,*

$$\left| \left\| \sum_{j=J'}^J \Lambda_{j,n} U_j(t_n) \right\|_{L^3}^3 - \sum_{j=J'}^J \left\| \Lambda_{j,n} U_j(t_n) \right\|_{L^3}^3 \right| < \epsilon$$

for all $n \geq N_0$.

To prove this lemma, using again the elementary inequality (2.25), we must estimate a finite number (depending on J) of terms of the form

$$\int |\Lambda_{j_1,n} U_{j_1}(t_n)| |\Lambda_{j_2,n} U_{j_2}(t_n)|^2$$

where $j_1, j_2 \in \{J', \dots, J\}$ and $j_1 \neq j_2$. Such a term tends to zero uniformly (for each fixed J) as $n \rightarrow \infty$ after passing to the following subsequence: if $t_n \lambda_{j_i,n}^{-2} \rightarrow \infty$ for $i = 1$ or $i = 2$, then necessarily the i th profile is global and using Hölder once more this term tends to zero by Theorem 8. If neither tends to infinity, we may pass to a subsequence so that the re-scaled times in both terms approach a constant time. Since we stay strictly away from the blow-up times of any profile, we may use the time continuity of solutions to replace the re-scaled time value by this fixed time in each profile, after which the orthogonality of the scales/cores implies that the term tends to zero due to (2.3). One may take a diagonal subsequence so that this is true in all instances when the times are bounded which proves Lemma 3.2. \square

3.3. Proof of Theorem 5. Consider the bounded sequence described in Theorem 5 and its profile decomposition after passing to a subsequence. For notational convenience, set $\lambda_{0,n} \equiv 1$ and $x_{0,n} \equiv 0$, so that the “0”th profile is the weak limit (without any transformations).

Note first that since $T^*(u_{0,n}) < +\infty$ for all n in the sequence which we consider, there must be at least one $j \geq 0$ such that $T^*(\phi_j) < \infty$.

If not, one could take $\tau_n \equiv \infty$ in Theorem 3 and hence u_n is global for large n , contrary to assumption.

Property (2.6) and small data results in $L^3(\mathbb{R}^3)$ now imply that there exists some $J_0 \in \mathbb{N}$ such that after re-ordering the profiles one has

$$T_j^* < \infty \iff 0 \leq j \leq J_0.$$

The orthogonality properties of the scales $\lambda_{j,n}$ then allow us to rearrange the first J_0 profiles in such a way that for all sufficiently large n , one has $\lambda_{0,n}^2 T_0^* \leq \lambda_{1,n}^2 T_1^* \leq \dots \leq \lambda_{J_0,n}^2 T_{J_0}^*$. Fix now such an ordering of the ϕ_j (so it is now possible that $\lambda_{0,n} \neq 1$, $x_{0,n} \neq 0$).

Note now that for any $s \in [0, T_0^*)$, we have

$$(3.1) \quad \Lambda_{0,n}^{-1} u_n(s) = U_0(s) + \sum_{j=1}^{J_0} \Lambda_{0,n}^{-1} \Lambda_{j,n} U_j(s) + \Lambda_{0,n}^{-1} w_n^J(s) + \Lambda_{0,n}^{-1} r_n^J(s).$$

We now claim that the left-hand side converges (after possibly passing to a subsequence) weakly in $L^3(\mathbb{R}^3)$ to $U_0(s)$, in which case by properties of weak limits and the invariance of the spatial norm we have

$$\|U_0(s)\|_{L^3(\mathbb{R}^3)} \leq \liminf_{n \rightarrow \infty} \|u_n(\lambda_{0,n}^2 s)\|_{L^3(\mathbb{R}^3)} \leq A_c$$

hence

$$\sup_{s \in [0, T_0^*)} \|U_0(s)\|_{L^3(\mathbb{R}^3)} \leq A_c.$$

On the other hand, since $T_0^* < \infty$, by definition of A_c we must also have $\sup_{s \in [0, T_0^*)} \|U_0(s)\|_{L^3(\mathbb{R}^3)} \geq A_c$ and hence U_0 is a critical element.

We shall now prove this weak convergence using the smallness of the remainders, the time-continuity of the evolution of the profiles and the decay of the global ones as well as the orthogonality of the scales/cores. To simplify notation, in the following we shall denote $\|f\|_p := \|f\|_{L_x^p}$.

Fix $\epsilon > 0$. We need to show that there exists a subsequence in n such that for any $\varphi \in \mathcal{C}_0^\infty(\mathbb{R}^3)$,

$$(3.2) \quad | \langle \lambda_{0,n} u_n(\lambda_{0,n} \cdot + x_{0,n}, \lambda_{0,n}^2 s) - U_0(s), \varphi \rangle | < \epsilon$$

for n sufficiently large, where $\langle \cdot, \cdot \rangle$ denotes the pairing between L^3 and $(L^3)' = L^{\frac{3}{2}}$, i.e. integration over \mathbb{R}^3 of the product. This will be accomplished by estimating the left-hand side of (3.2) by

$$(3.3) \quad \sum_{j=1}^{J_1} | \langle \Lambda_{j,n} \Lambda_{0,n}^{-1} U_j(s), \varphi \rangle | + \left\| \sum_{j=J_1+1}^{J_0} \Lambda_{j,n} \Lambda_{0,n}^{-1} U_j(s) \right\|_3 \|\varphi\|_{\frac{3}{2}} + \langle \Lambda_{0,n}^{-1} w_n^J(s) + \Lambda_{0,n}^{-1} r_n^J(s), \varphi \rangle$$

where J_1 will be some fixed appropriately large integer. According to (2.6), J_1 can be chosen so large that for any $j > J_1$, one has

$$\|U_j(0)\|_3 = \|\phi_j\|_3 \leq \varepsilon_0$$

and $U_j = NS(\phi_j)$ can be produced by a fixed-point argument on $(0, \infty)$ (see, e.g., [9]). Moreover,

$$(3.4) \quad \sup_{t \geq 0} \|NS(\phi_j)(t)\|_3 \lesssim \|\phi_j\|_3$$

as a by-product of the small data theory. Hence, for such a J_1 we have

$$(3.5) \quad \sum_{j=J_1+1}^J \|\Lambda_{j,n} \Lambda_{0,n}^{-1} U_j(s)\|_3^3 \leq C_0 \sum_{j=J_1+1}^J \|\phi_j\|_3^3$$

for some universal $C_0 > 0$. Moreover, by up to a harmless rescaling by $\Lambda_{0,n}^{-1}$ in its statement (as it does not change the orthogonality of scales/cores), we may use Lemma 3.2 and pass to a subsequence in n so that for each J and any $J' < J$, there exists $n_0(J)$ (in fact, independent of J') such that for $n \geq n_0(J)$ one has

$$(3.6) \quad \left\| \sum_{j=J'}^J \Lambda_{j,n} \Lambda_{0,n}^{-1} U_j(s) \right\|_3^3 \leq 2 \sum_{j=J'}^J \|\Lambda_{j,n} \Lambda_{0,n}^{-1} U_j(s)\|_3^3.$$

Now, due to (2.6), by taking J_1 large enough (depending on φ), setting $J' = J_1 + 1$ in (3.6) and using the previous estimate we can make the middle term of (3.3) less than $\epsilon/3$ whenever $n \geq n_0(J)$ for $J > J_1$. Now, by property (2.14) and the scaling of the norm, there exists $n(J) \geq n_0(J)$, increasing in J , such that

$$\|\Lambda_{0,n(J)}^{-1} w_{n(J)}^J(s)\|_{\dot{B}_{p,p}^{s_p}}, \|\Lambda_{0,n(J)}^{-1} r_{n(J)}^J(s)\|_{\dot{B}_{p,p}^{s_p}} \rightarrow 0$$

as $J \rightarrow \infty$. In particular, these limits hold weakly. By heat estimates, (2.7) and the transformational invariance of the norm, $\Lambda_{0,n}^{-1} w_n^J(s)$ is bounded in L^3 . We therefore also have $\Lambda_{0,n}^{-1} r_n^J(s)$ bounded in L^3 by (3.1), our assumption that $u_n(t)$ is uniformly bounded in L^3 for all times and (3.6) with $J' = 0$ in conjunction with (2.6) to bound the other terms. Therefore the error terms tend weakly to zero in L^3 as well, and hence, setting $n = n(J)$, the third term in (3.3) can be made less than $\epsilon/3$ for sufficiently large J .

Finally, since $n(J) \rightarrow \infty$ with J , the orthogonality of the scales/cores shows that each term in the sum on the left in (3.3) tends to zero after a subsequence, by arguments similar to those in the proof of Lemma 3.2. Since there are only a finite number J_1 of these, the first term in (3.3) can be made less than $\epsilon/3$ for sufficiently large J which proves (3.2) for some subsequence of n 's as desired. \square

3.4. Proof of Theorem 6. Suppose now $u = NS(u_0)$ is a critical element, and consider the bounded sequence $u_{0,n} := u(s_n)$ for some $s_n \nearrow T^*(u_0)$. Pass to a subsequence so that one may write $u_{0,n}$ and $u_n := NS(u_{0,n})$ in terms of profiles with the notations of Theorems 1 and 3.

As in the proof of Theorem 5 above, there is some $J_0 \geq 0$ such that $T_j^* < \infty \iff 0 \leq j \leq J_0$ and we may re-arrange the first J_0 profiles in such a way that $\lambda_{0,n}^2 T_0^* \leq \lambda_{1,n}^2 T_1^* \leq \dots \leq \lambda_{J_0,n}^2 T_{J_0}^*$ for all sufficiently large n . Fix now such an ordering of the ϕ_j , and suppose that 0 has been moved now to some $j_0 \in \mathbb{N}$, that is, now $\lambda_{j_0,n} \equiv 1$ and $x_{j_0,n} \equiv 0$ and ϕ_{j_0} is the weak limit of $u_{0,n}$. The theorem will therefore be proved if we show that $\phi_{j_0} = 0$.

By the definition of τ_n , etc., in Theorem 3, it is clear that with this ordering we must have

$$(3.7) \quad \lambda_{0,n}^2 T_0^* \leq T^*(u_{0,n}) = T^*(u_0) - s_n$$

for large n , and hence $\lambda_{0,n} \rightarrow 0$ as $n \rightarrow \infty$. In particular, we see that $j_0 \neq 0$, that is, 1 (the scale of the weak limit profile) cannot be smaller than all other scales.

We shall need the following crucial claim (which actually applies to the more general sequence $\{u_{0,n}\}$ considered in the proof of Theorem 5), whose proof we postpone momentarily:

Claim 3.3. *Fix any $s \in (0, T_0^*)$. Setting $t_n := \lambda_{0,n}^2 s$, after possibly passing to a subsequence in n one has*

$$\|u_n(t_n)\|_{L^3}^3 \geq \|\Lambda_{0,n} U_0(t_n)\|_{L^3}^3 + \|u_n(t_n) - \Lambda_{0,n} U_0(t_n)\|_{L^3}^3 + o(1)$$

as $n \rightarrow \infty$.

Let us proceed to prove Theorem 6. Exactly as in the proof of Theorem 5, we see again that U_0 is a critical element since we have

$$A_n := \sup_{t \in [0, T^*(u_{0,n})]} \|NS(u_{0,n})(t)\|_{L^3} = \sup_{t \in [s_n, T^*(u_0)]} \|u(t)\|_{L^3} \equiv A_c$$

for all n , due to the definition of A_c and the fact that $T^*(u_0) < \infty$. We shall now show that this implies by Claim 3.3 that $\phi_{j_0} = 0$, i.e. that $u_{0,n}$ tends weakly to zero which was our goal. Fix any $\epsilon > 0$. By the time-continuity of solutions we may take $s \in (0, T_0^*)$ such that

$$A_c^3 - \|U_0(s)\|_3^3 < (\epsilon/2)^3 C_0^{-3}$$

where $C_0 > 0$ is the universal constant in the continuous embedding $L^3 \hookrightarrow \dot{B}_{p,p}^{s_p}$. Set $t_n := \lambda_{0,n}^2 s$. Then due to Claim 3.3, after passing

to a further subsequence in n , we have

$$\begin{aligned} A_c^3 &\geq \|u_n(t_n)\|_3^3 \geq \|U_0(s)\|_3^3 \\ &\quad + C_0^{-3} \left\| \sum_{j=1}^J \Lambda_{j,n} U_j(t_n) + w_n^J(t_n) + r_n^J(t_n) \right\|_{\dot{B}_{p,p}^{s_p}}^3 + C_0^{-3} \epsilon(n, s) \end{aligned}$$

for any J where $\epsilon(n, s) \rightarrow 0$ as $n \rightarrow \infty$. According to (2.14), we may fix $J \geq j_0$ so large that

$$\|w_n^J(t_n) + r_n^J(t_n)\|_{\dot{B}_{p,p}^{s_p}} \leq \epsilon/2$$

for sufficiently large n . The previous two inequalities give

$$\begin{aligned} ((\epsilon/2)^3 - \epsilon(n, s))^{\frac{1}{3}} + \epsilon/2 &\geq \left\| \sum_{j=1}^J \Lambda_{j,n} U_j(t_n) \right\|_{\dot{B}_{p,p}^{s_p}} \\ &=: \left(\sum_{j=1}^J \|\Lambda_{j,n} U_j(t_n)\|_{\dot{B}_{p,p}^{s_p}}^p - \epsilon_J(n) \right)^{\frac{1}{p}}, \end{aligned}$$

and we claim (as in Lemma 3.2) that $\epsilon_J(n) \rightarrow 0$ as $n \rightarrow \infty$ for fixed J , after passing to a subsequence in n . Postponing this fact for a moment, we have now shown that

$$\|U_{j_0}(t_n/\lambda_{j_0,n}^2)\|_{\dot{B}_{p,p}^{s_p}}^p \leq \left(((\epsilon/2)^3 - \epsilon(n, s))^{\frac{1}{3}} + \epsilon/2 \right)^p + \epsilon_J(n).$$

Recall that $\lambda_{j_0,n} \equiv 1$ because φ_{j_0} is the weak limit of $u_{0,n}$, and note that $t_n = \lambda_{0,n}^2 s \rightarrow 0$ as $n \rightarrow \infty$ due to (3.7). Therefore letting $n \rightarrow \infty$ and using the continuity of U_{j_0} in $\dot{B}_{p,p}^{s_p}$ at $t = 0$ we have

$$\|U_{j_0}(0)\|_{\dot{B}_{p,p}^{s_p}} = \|\varphi_{j_0}\|_{\dot{B}_{p,p}^{s_p}} \leq \epsilon.$$

Since ϵ was arbitrary and $\phi_{j_0} \in L^3$, this implies that $\phi_{j_0} = 0$, which proves the theorem.

All that remains now is to show that $\epsilon_J(n) \rightarrow 0$, which we now explain. It is again a simple consequence of the orthogonality of the scales/cores, and is proved in the same way as Lemmas 3.2 and 2.6. In fact, up to undoing a harmless $\Lambda_{0,n}^{-1}$ transform, we follow closely the proof of Lemma 2.6, without the inner L_t^t norm and with different times in the profiles U_i : using the elementary inequality (2.25), in the

expansion of the Besov norm of the sum $\sum_{j=j_0+1}^J U_n^{j,0}$, we get $O(J)$ cross-terms like

$$\begin{aligned} & \int_0^\infty \tau^\gamma \int \left| K(\tau) \left[\frac{1}{\lambda_{1,n}} U_1 \left(\frac{\cdot - x_{1,n}}{\lambda_{1,n}}, s_1 \right) \right] (x) \right| \\ & \quad \times \left| K(\tau) \left[\frac{1}{\lambda_{2,n}} U_2 \left(\frac{\cdot - x_{2,n}}{\lambda_{2,n}}, s_2 \right) \right] (x) \right|^{p-1} dx d\tau = \\ & \int_0^\infty \tau^\gamma \int \left| \frac{1}{\lambda_{1,n}} [K(\lambda_{1,n}^{-2}\tau)[U_1(s_1)]] \left(\frac{x - x_{1,n}}{\lambda_{1,n}} \right) \right| \\ & \quad \times \left| \frac{1}{\lambda_{2,n}} [K(\lambda_{2,n}^{-2}\tau)[U_2(s_2)]] \left(\frac{x - x_{2,n}}{\lambda_{2,n}} \right) \right|^{p-1} dx d\tau \end{aligned}$$

for some $s_1, s_2 > 0$ in the life-spans of U_1 and U_2 respectively (here we have passed to a subsequence and used the time continuity of the profiles and Theorem 8 as in the proof of Lemma 3.2). Using the change of variables $y = \frac{x - x_{1,n}}{\lambda_{1,n}}$, $s = \lambda_{1,n}^{-2}\tau$, we see that this term equals

$$\left(\frac{\lambda_{1,n}}{\lambda_{2,n}} \right)^{\frac{2+d}{p'}} \int_0^\infty \int |V_1(y, s)| \left| V_2 \left(\frac{\lambda_{1,n}}{\lambda_{2,n}} y + \frac{x_{1,n} - x_{2,n}}{\lambda_{2,n}}, \left(\frac{\lambda_{1,n}}{\lambda_{2,n}} \right)^2 s \right) \right|^{p-1} dy ds,$$

where $\frac{1}{p} + \frac{1}{p'} = 1$ and $V_i(y, s) = s^{\frac{\gamma}{p}} [K(s)[U_i(s_i)]](y)$ for $i \in \{1, 2\}$. Since $U_i(s_i) \in \dot{B}_{p,p}^{s_p}$, by (2.23) we can approximate V_1 and V_2 in $L^p(\mathbb{R}^d \times (0, +\infty))$ by smooth functions of (y, s) with compact support in $\mathbb{R}^d \times (0, +\infty)$. It is therefore clear that the above term tends to zero as $n \rightarrow \infty$ if $\frac{\lambda_{1,n}}{\lambda_{2,n}} \rightarrow 0$ as $n \rightarrow \infty$, or if $\lambda_{1,n} \equiv \lambda_{2,n}$ and $\left| \frac{x_{1,n} - x_{2,n}}{\lambda_{2,n}} \right| \rightarrow \infty$ as $n \rightarrow \infty$. Similarly, if we appropriately chose the new variables in terms of $\lambda_{2,n}$ and $x_{2,n}$ instead, we can easily show that the term also tends to zero if $\frac{\lambda_{2,n}}{\lambda_{1,n}} \rightarrow 0$, and the result follows in view of (2.1). \square

Remark 3.4. *A similar argument can be used to show that only one profile can be a critical element since all others are small at some time, implying that they exist globally by the small data theory. Although this fact was used to prove the theorems in [14], we shall not use it here.*

Remark 3.5. *One could also prove a similar compactness result as Theorem 3.2 in [14], namely that if $NS(u_0)$ satisfies*

$$\sup_{t \in [0, T^*(u_0))} \|NS(u_0)(t)\|_{L^3} = A_c$$

then for any sequence $\{t_n\}$ converging to $T^*(u_0)$, there exists a sequence $\{s_n\}$ with $t_n \leq s_n \nearrow T^*(u_0)$ such that sequence $NS(u_0)(s_n)$ is compact in $\dot{B}_{p,p}^{sp}$ up to norm-invariant transformations in space.

Remark 3.6. Claim 3.3 also immediately proves Theorem 5, but we feel that the proof given above is more self-contained and perhaps more natural at that point.

Proof of Claim 3.3. Note first that we may assume without loss of generality that u_n is scalar-valued by setting

$$\|(f^k)_{k=1}^3\|_{L^3} := \|(\|f^k\|_{L^3})_{k=1}^3\|_{\ell^3}$$

and treating each component separately.

We first remark that, after passing to an appropriate subsequence,

$$(3.8) \quad C_1 := \sup_{\substack{J \geq 0 \\ n \geq N_0(J)}} \left\| \sum_{j=0}^J \Lambda_{j,n} U_j(t_n) \right\|_3 < \infty \quad \text{and}$$

$$C_2 := \sup_{\substack{J \geq 0 \\ n \geq N_0(J)}} \|R_n^J(t_n)\|_3 < \infty ,$$

where $R_n^J = w_n^J + r_n^J$ and $N_0(J)$ is as in Lemma 3.2 with $\epsilon = \epsilon(J)$ chosen appropriately. Indeed, C_2 is bounded by C_1 and $\sup_n A_n$ (where we recall that $A_n = \sup_{0 \leq t < T^*(u_{0,n})} \|u_n(t)\|_3$ and $A_n \rightarrow A_c < \infty$). To show

$C_1 < \infty$, for J large and $J_1 < J$ we can write

$$\left\| \sum_{j=0}^J \Lambda_{j,n} U_j(t_n) \right\|_3 \leq \left\| \sum_{j=0}^{J_1} \Lambda_{j,n} U_j(t_n) \right\|_3 + \left\| \sum_{j=J_1+1}^J \Lambda_{j,n} U_j(t_n) \right\|_3 .$$

For J_1 sufficiently large, (2.6), Lemma 3.2 and (3.4) give a uniform bound of the second term. Since we stay strictly away from the potential blow-up times of all profiles, for fixed J_1 the first term is bounded due to Theorem 8 and the time-continuity in L^3 of each U_j .

Set $v_n := u_n - \Lambda_{0,n} U_0$. Due to (2.25), we have

$$\begin{aligned} \left| \|u_n(t_n)\|_3^3 - \|\Lambda_{0,n} U_0(t_n)\|_3^3 - \|v_n(t_n)\|_3^3 \right| &\lesssim \int |\Lambda_{0,n} U_0(t_n)|^2 |v_n(t_n)| \\ &\quad + \int |\Lambda_{0,n} U_0(t_n)| |v_n(t_n)|^2 . \end{aligned}$$

Unlike in [14] where there is only a cross-term similar to the first one (hence one may leave the absolute value outside the integral and use weak convergence), both terms require the use of specific information about the components of v_n . We deal with the second term first and

then briefly indicate how the first one can be dealt with in a similar way.

We would therefore like to show that

$$\int |\Lambda_{0,n}U_0(t_n)||v_n(t_n)|^2 \rightarrow 0$$

as $n \rightarrow \infty$ along some subsequence. Fix some $L_0 \in \mathbb{N}$ large, to be chosen precisely later. Then for $J > L_0$ we write

$$v_n(t_n) = \sum_{j=1}^{L_0} \Lambda_{j,n}U_j(t_n) + \sum_{j=L_0+1}^J \Lambda_{j,n}U_j(t_n) + R_n^J(t_n) =: A_1 + A_2 + A_3$$

so that

$$(v_n(t_n))^2 = (A_1)^2 + (A_2)^2 + (A_3)^2 + 2(A_1A_2 + A_2A_3 + A_1A_3) .$$

For the first term, note that using the arguments in the proof of Lemma 3.2 we can make $\int |\Lambda_{0,n}U_0(t_n)||A_1|^2$ arbitrarily small for sufficiently large n (depending on L_0 , which we shall fix in a moment) by orthogonality of the scales/cores. For the second term, using Hölder's inequality and Lemma 3.2, for an appropriate subsequence of n depending on J , we have

$$\int |\Lambda_{0,n}U_0(t_n)||A_2|^2 \leq \|U_0\|_3 \left(2 \sum_{j=L_0+1}^J \|\Lambda_{j,n}U_j(t_n)\|_3^3 \right)^{\frac{2}{3}}$$

which can be made arbitrarily small by choosing L_0 sufficiently large due to (2.6) and arguments similar to the proof of (3.5). Using arguments similar to those above and noting that $\|A_1\|_3$ and $\|A_3\|_3$ are uniformly bounded by (3.8), we can treat all remaining terms except for those of the form

$$\int |\Lambda_{0,n}U_0(t_n)||w_n^J(t_n)|^2 \quad \text{or} \quad \int |\Lambda_{0,n}U_0(t_n)||r_n^J(t_n)|^2 .$$

Since $U_0(s) \in L^3$, using Hölder's inequality, (2.7) and heat estimates we can control the term involving w_n^J by a uniform constant times the quantity

$$\|\Lambda_{0,n}U_0(t_n)w_n^J(t_n)\|_{\frac{3}{2}} .$$

Then recalling that $t_n = \lambda_{0,n}^2 s$ and $s_p = -1 + \frac{3}{p}$, approximating $U_0(s)$ in L^3 by a smooth compactly supported function and noting that one may replace $K(\tau)$ by $e^{\tau\Delta}$ in (2.23) to obtain yet another equivalent

Besov norm, we can control this term by

$$\begin{aligned} \|\Lambda_{0,n}U_0(t_n)\|_{p'}\|w_n^J(t_n)\|_p &= \|U_0(s)\|_{p'}\lambda_{0,n}^{-s_p}\|e^{tn\Delta}\psi_n^J\|_p \\ &\lesssim s^{-\frac{s_p}{2}}\|U_0(s)\|_{p'}\|\psi_n^J\|_{\dot{B}_{p,\infty}^{s_p}} \end{aligned}$$

where $\frac{2}{3} = \frac{1}{p'} + \frac{1}{p}$. As s is fixed, this term can therefore be made small for large J and then n due to (2.9) and the continuous embedding $\dot{B}_{p,p}^{s_p} \hookrightarrow \dot{B}_{p,\infty}^{s_p}$.

We now just need to show that $\int |\Lambda_{0,n}U_0(t_n)||r_n^J|^2$ can be made arbitrarily small for large J and n . By a change of variables, we have

$$\int |\Lambda_{0,n}U_0(x, t_n)||r_n^J(x, t_n)|^2 dx = \int |U_0(y, s)| |\Lambda_{0,n}^{-1}r_n^J(y, s)|^2 dy .$$

Set $\tilde{r}_n^J := \Lambda_{0,n}^{-1}r_n^J(s)$. Note that $\sup_{n,J} \|\tilde{r}_n^J\|_{L^{\frac{3}{2}}} < \infty$ by (3.8), (2.7) and linear heat estimates, and we may assume that $U_0 \in \mathcal{C}_0^\infty$ by approximation in L^3 . It therefore suffices to show that $|\tilde{r}_{n(J)}^J(y)|^2$ tends strongly to zero (for some increasing $n(J)$) as $J \rightarrow \infty$ in some Banach space $B \hookrightarrow \mathcal{S}'$ and hence tends weakly to zero in $L^{\frac{3}{2}}$, making this term small for large J and $n = n(J)$.

In order to do this, we claim that, since r_n^J satisfies an equation of the form (2.18), there exists $N(J) \in \mathbb{N}$ defined for all $J \geq 0$ such that

$$(3.9) \quad \sup_{\substack{J \geq 0 \\ n \geq N(J)}} \|\tilde{r}_n^J\|_{\dot{B}_{3/2,\infty}^1} < \infty .$$

Let us postpone the proof of (3.9) for a moment, and use it to complete the proof of Claim 3.3.

Recalling standard product estimates in Besov spaces, we have

$$\|fg\|_{\dot{B}_{3/2,\infty}^\sigma} \lesssim \|f\|_{\dot{B}_{p,\infty}^{s_p}} \|g\|_{\dot{B}_{3/2,\infty}^{\sigma+1}}$$

for any $\sigma > 0$. Note that this is a valid application of the product laws since $s_p + (\sigma + 1) = \frac{3}{p} + \sigma > 0$ and $s_p < \frac{3}{p}$. Therefore we have

$$\|(\tilde{r}_n^J)^2\|_{\dot{B}_{3/2,\infty}^0} \lesssim \|\tilde{r}_n^J\|_{\dot{B}_{p,\infty}^{s_p}} \|\tilde{r}_n^J\|_{\dot{B}_{3/2,\infty}^1} \lesssim \|\tilde{r}_n^J\|_{\dot{B}_{p,p}^{s_p}} \|\tilde{r}_n^J\|_{\dot{B}_{3/2,\infty}^1} .$$

Hence (2.14) along with (3.9) imply that

$$\| |\tilde{r}_{n(J)}^J|^2 \|_{\dot{B}_{3/2,\infty}^0} \rightarrow 0 \quad \text{as } J \rightarrow \infty$$

for some $n(J)$ increasing in J , which concludes the proof.

We now briefly return to the proof of (3.9), which is nothing but a simple consequence of estimates on the Duhamel term in [6]. Indeed,

the proof of Proposition 1 in [6] gives the estimate

$$(3.10) \quad \|B(f, g)(t)\|_{\dot{B}_{3/2, \infty}^1} \lesssim \sup_{0 < \tau < t} \|fg(\tau)\|_{L^{3/2}}$$

where $B(f, g) = (\partial_t - \Delta)^{-1} \mathbb{P} \nabla \cdot (f \otimes g)$ with $B(f, g)(0) = 0$. According to (2.18) we can write r_n^J as a sum of a finite number (independent of J) of terms (each of which do depend on J) of the form $B(f, g)$ which can all be controlled, after applying (3.10) and then Hölder in x , by (3.8), (2.7) and standard heat estimates, plus a sum of the form

$$\sum_{0 \leq j \neq j' \leq J} B(\Lambda_{j, n} U_j, \Lambda_{j', n} U_{j'}) .$$

After applying (3.10), we can bound this term by a constant independent of J by the orthogonality of the scales/cores (as in the proof of Lemma 3.2) for $n \geq N(J)$ sufficiently large for any J . Applying $\Lambda_{0, n}^{-1}$ (under which all norms concerned are invariant) everywhere establishes (3.9) and we are done with the quadratic term $|v_n|^2$.

We now go back to proving

$$\int |\Lambda_{0, n} U_0(t_n)|^2 |v_n(t_n)| \rightarrow 0 ,$$

using the same decomposition of v_n as a sum of three terms. Applying the triangle inequality, terms with A_1 and A_2 go to zero by the same arguments of orthogonality of scales/cores. Similarly, the term in A_3 involving w_n^J goes to zero using the heat decay estimates. Hence all we are left with is

$$\int |U_0(y, s)|^2 |\Lambda_{0, n}^{-1} r_n^J(y, s)| \, dy .$$

We just proved that $r_n^J \in \dot{B}_{3/2, \infty}^1$, while we know that r_n^J goes to zero in, say, $\dot{B}_{4, 4}^{-1/4}$, where we take $p = 4$ for concreteness (general p is treated similarly). By interpolation, we get that $r_n^J \in \dot{B}_{2, 20/3}^{1/2}$ and goes to zero in this later norm. We conclude using composition rules in Besov spaces $\dot{B}_{p, q}^s$, with $0 < s < 1$, as $\| |f| \|_{\dot{B}_{p, q}^s} \lesssim \|f\|_{\dot{B}_{p, q}^s}$ for such s (a fact which readily follows from the characterization of Besov spaces in terms of finite differences in that range, and the elementary inequality $\||a| - |b|\| \leq |a - b|$). As $|U_0|^2$ is smooth, hence in the dual space $\dot{B}_{2, 20/17}^{-1/2}$, this last remaining integral goes to zero, and Claim 3.3 is proved. \square

3.5. Proof of Theorem 7. Theorem 7 is a consequence of the following lemma which is proved in the last section of [14], following the argument in [8]:

Lemma 3.7. *Suppose $u_0 \in L^3$ and $NS(u_0)$ belongs to $L^\infty([0, T]; L^3)$ for some finite $T > 0$. Then there exists some $R_0 \gg 1$ such that u belongs to $C^\infty((\mathbb{R}^3 \setminus B_{R_0}(0)) \times [0, T])$, with global bounds on derivatives.*

Indeed, assuming $T^*(u_0) < +\infty$ and applying Lemma 3.7 with $T = T^*(u_0)$, $u(t) \rightarrow 0$ as $t \nearrow T^*(u_0)$ implies that $D^\alpha u(x, T^*) \equiv 0$ for $|x| > R_0$ for any multi-index α . Then known backwards uniqueness and unique continuation results for the parabolic inequality satisfied by $\omega := \nabla \times u$ show that $\omega \equiv 0$ on $\mathbb{R}^3 \times [0, T^*]$, see the last section of [14] for more details (see also [8]). This implies $u \equiv 0$ as well due to the divergence-free condition, and hence $T^*(u_0) = +\infty$ by uniqueness of mild solutions, contrary to assumption, which proves Theorem 7. \square

4. MINIMAL BLOW-UP INITIAL DATA

In this section we consider the question of “minimal blow-up initial data” in various settings, of the type addressed in [23].

Suppose $X = X(\mathbb{R}^d)$ is a Banach space of initial data on which there is a norm which is invariant under the transformations leaving the Navier-Stokes equations invariant, with the property that there exists some small $\epsilon_0 = \epsilon_0(X) > 0$ such that $T^*(u_0) = +\infty$ whenever $\|u_0\|_X < \epsilon_0$, where $T^*(u_0)$ is the maximal time of existence of $NS(u_0)$ in the space X .

Then the question to be considered is the following:

Statement 4.1. *Suppose there exists $v_0 \in X$ such that $T^*(v_0) < \infty$, and define $\rho = \rho_X := \inf\{\|v_0\|_X \mid T^*(v_0) < +\infty\} \geq \epsilon_0 > 0$. Then there exists $u_0 \in X$ such that $T^*(u_0) < \infty$, and $\|u_0\|_X = \rho$. Moreover, up to transformations under which the Navier-Stokes equations are invariant, the set of such u_0 is compact in Y , for a similar space Y such that $X \subseteq Y$.*

Such a statement was proved in the case $X = \dot{H}^{\frac{1}{2}}(\mathbb{R}^3)$ in [23] (in fact in the setting of weak solutions), and moreover with $Y = X$. In the following we show that the result is a simple consequence of the profile decompositions, Theorem 2 in [10] for $X = \dot{H}^{\frac{1}{2}}$ (and in fact this can easily be extended to $X = \dot{H}^{\frac{d}{2}-1}(\mathbb{R}^d)$ for any d), Theorems 1 and 2 stated above (proved in [19]) and Theorem 3 stated and proved above in the settings $X = L^d(\mathbb{R}^d)$ and $X = \dot{B}_{a,b}^{s_a}(\mathbb{R}^d)$. To be precise, what we prove is the following:

Theorem 9. *Statement 4.1 is true for $X = Y = \dot{H}^{\frac{d}{2}-1}(\mathbb{R}^d)$ for any $d \geq 2$, and there exists a norm on X , equivalent to the standard norm, such that Statement 4.1 is true for $(X, Y) = (L^d(\mathbb{R}^d), \dot{B}_{p,q}^{s_p}(\mathbb{R}^d))$*

whenever $3 \leq d < p \leq q \leq \infty$, and for $(X, Y) = (\dot{B}_{a,b}^{s_a}(\mathbb{R}^d), \dot{B}_{p,q}^{s_p}(\mathbb{R}^d))$ for any $d \geq 3$ and $a, b \in [1, 2d + 3)$ satisfying $\max\{a, b\} < p \leq \infty$ and $1 \leq b < (p/a)b \leq q \leq \infty$, where $s_r := -1 + \frac{d}{r}$ for $r \in \mathbb{R}$.

Note that it is important that $b < \infty$ so that local solutions are in fact known to exist (and hence a maximal time of existence makes sense), as opposed to only having global solutions for small data. (We shall see below that this is necessary for a different technical reason as well.) Note also that in applying Theorem 3 in the proof below, one may have to use a set of smaller indices first in the space Y (to satisfy the assumptions of that theorem), and then the more general statement follows from the standard embeddings.

Proof of Theorem 9. For simplicity, we first prove the theorem for $X = \dot{B}_{a,b}^{s_a}$, and define $\|\cdot\|_X := \|\cdot\|_{\dot{B}_{a,b}^{s_a}}$ (this norm is defined via wavelet bases, see [19]). Assume there is some $\dot{B}_{a,b}^{s_a}$ datum with a finite maximal time of existence, so that $\rho = \rho_{\dot{B}_{a,b}^{s_a}}$ is well-defined. By known small data regularity results there exists ϵ_0 such that $\rho \geq \epsilon_0 > 0$. By the definition of ρ , there exists a sequence $u_{0,n} \in \dot{B}_{a,b}^{s_a}$ with $T^*(u_{0,n}) < \infty$ (hence necessarily $\|u_{0,n}\|_{\dot{B}_{a,b}^{s_a}} \geq \rho$) and $\|u_{0,n}\|_{\dot{B}_{a,b}^{s_a}} \searrow \rho$ as $n \rightarrow \infty$.

Since $u_{0,n}$ is therefore a bounded sequence in $\dot{B}_{a,b}^{s_a}$, we can apply the profile decomposition Theorem 2 to this sequence, so that, after passing to a subsequence, we may write $u_{0,n}$ as

$$u_{0,n}(x) = \sum_{j=0}^J \frac{1}{\lambda_{j,n}} \phi_j \left(\frac{x - x_{j,n}}{\lambda_{j,n}} \right) + \psi_n^J(x),$$

and (2.10) gives

$$(4.1) \quad \sum_{j=0}^{\infty} \|\phi_j\|_{\dot{B}_{a,b}^{s_a}}^{\tau} \leq \liminf_{n' \rightarrow \infty} \|u_{0,n'}\|_{\dot{B}_{a,b}^{s_a}}^{\tau} = \rho^{\tau}.$$

Moreover, applying Theorem 3, we see that there is at least one $j_0 \in \mathbb{N}$ such that $T^*(\phi_{j_0}) < +\infty$. Indeed, otherwise one could take $\tau_n \equiv +\infty$ in that theorem and see that $NS(u_{0,n})$ lives past its finite maximal time of existence which is impossible. By definition of ρ , we know that $\|\phi_{j_0}\|_{\dot{B}_{a,b}^{s_a}} \geq \rho$, else we would have $T^*(\phi_{j_0}) = +\infty$. However, (4.1) gives $\|\phi_{j_0}\|_{\dot{B}_{a,b}^{s_a}} \leq \rho$, so that $\|\phi_{j_0}\|_{\dot{B}_{a,b}^{s_a}} = \rho$ and we may take $u_0 = \phi_{j_0}$ in the statement of the Theorem. This proves the existence statement.

Remark 4.2. We remark similarly to Remark 3.4 in the previous section that (4.1) implies that only one profile appears in the decomposition of $u_{0,n}$.

To prove the compactness statement, suppose now moreover that

$$\|u_{0,n}\|_{\dot{B}_{a,b}^{s_a}} \equiv \rho$$

(having established the existence of at least one such element) and passing to a subsequence write $u_{0,n}$ in a profile decomposition as before. The same results hold as above, and note that as pointed out in Remark 4.2, (4.1) implies that $\phi_j = 0$ for all $j \neq j_0$, since necessarily $\|\phi_{j_0}\|_{\dot{B}_{a,b}^{s_a}} = \rho$ implies that $\sum_{j \neq j_0} \|\phi_j\|_{\dot{B}_{a,b}^{s_a}}^\tau \leq 0$. Note that here we have used the fact that $b < \infty$ so that $\tau < \infty$. Therefore we can write

$$u_{0,n}(x) = \frac{1}{\lambda_n} \phi \left(\frac{x - x_n}{\lambda_n} \right) + \psi_n(x) =: \Lambda_n \phi(x) + \psi_n(x),$$

where $\phi = \phi_{j_0}$, etc., and $\psi_n \rightarrow 0$ in $\dot{B}_{p,q}^{s_p}$ as $n \rightarrow \infty$ by (2.9). The invariances of the $\dot{B}_{p,q}^{s_p}$ norm imply as well that $\Lambda_n^{-1} \psi_n \rightarrow 0$ in $\dot{B}_{p,q}^{s_p}$, hence clearly $\Lambda_n^{-1} u_{0,n} \rightarrow \phi$ in $\dot{B}_{p,q}^{s_p}$, and the theorem is proved.

To prove the theorem for $X = L^d(\mathbb{R}^d)$, we consider a minimizing sequence in $L^d(\mathbb{R}^d)$ and proceed as above applying Theorem 3. We similarly conclude that there exists some profile ϕ_{j_0} with finite maximal time of existence in $\dot{B}_{a,b}^{s_a}$. Otherwise, $NS(u_{0,n})$ would be globally defined in $\dot{B}_{a,b}^{s_a}$, and standard ‘‘persistence’’ results for Navier-Stokes (see, e.g., [9]) would then imply that $NS(u_{0,n})$ is global in $L^d(\mathbb{R}^d)$ as well, contrary to assumption. The remainder of the proof follows as above due to Remark 2.3 and the theorem is proved in this case as well.

For $X = \dot{H}^{\frac{d}{2}-1}(\mathbb{R}^d)$, the proof is identical using the theorems in [10] (with the usual norm on X) and we would initially take $Y = L^d(\mathbb{R}^d)$ to see that $\Lambda_n^{-1} u_{0,n} \rightarrow \phi$ in Y . Moreover, since $X \hookrightarrow Y$, $\phi \in X$ and $\|\Lambda_n^{-1} u_{0,n}\|_X = \|u_{0,n}\|_X \equiv \|\phi\|_X = \rho$, we see that $\Lambda_n^{-1} u_{0,n} \rightharpoonup \phi$ in X , and since X is a Hilbert space the above properties imply that actually we have strong convergence in X as desired. \square

APPENDIX A. A PERTURBATION RESULT

Let us state (without proof) a perturbation result for the d -dimensional Navier-Stokes system.

Proposition A.1. *Let $s_p = -1 + \frac{d}{p}$, $r \in [1, \frac{2p}{p+1}]$ and define $s := s_p + \frac{2}{r}$. Assume finally that $p < 2d + 3$. There are constants ε_0 and C such that the following holds. Let $w_0 \in \dot{B}_{p,p}^{s_p}$, $f \in F := \mathcal{L}^r([0, T]; \dot{B}_{p,p}^{s-2}) + \mathcal{L}^{\frac{2p}{p+1}}([0, T]; \dot{B}_{p,q}^{s_p-1+\frac{1}{p}})$ and $v \in D := \mathcal{L}^p([0, T]; \dot{B}_{p,p}^{s_p+\frac{2}{p}})$ be given, with*

$$\|w_0\|_{\dot{B}_{p,p}^{s_p}} + \|f\|_F \leq \varepsilon_0 \exp(-C\|v\|_D).$$

Suppose moreover that $\operatorname{div} v = 0$, and let w be a solution of

$$\partial_t w - \Delta w + w \cdot \nabla w + v \cdot \nabla w + w \cdot \nabla v = -\nabla \pi + f$$

with $\operatorname{div} w = 0$. Then w belongs to $E_{p,p}(T)$ and the following estimate holds:

$$\|w\|_{E_{p,p}(T)} \lesssim (\|w_0\|_{\dot{B}_{p,q}^{s_p}} + \|f\|_F) \exp C\|v\|_D.$$

The proof of that proposition follows the estimates of [9] (see in particular Propositions 4.1 and Theorem 3.1 of [9]). The two main differences are

- the absence of an exterior force in [9], but that force is added with no difficulty to the estimates ;
- the rather weak estimate on the drift term v , which accounts for the restricted numerology on time exponents in the definition of $E_{p,p}$. The reader should note that closing estimates on w in our setting amounts to doing again the same estimates that were done in the proof of Lemma 2.7.

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