

Fourier-Whittaker expansions:  
from analysis on global fields to  
geometry on local fields

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# ① Fourier-Whittaker expansions: global theory

An important tool for the study of automorphic forms and their L-functions is their Fourier expansion / Whittaker coefficients.

Let's recall this briefly for  $GL_2$  (would work similarly for  $GL_n$ ). Let  $E =$  global field, let  $f$  be an automorphic form for  $GL_2/E$ .

For any  $g \in GL_2(\mathbb{A}_E)$ , can write the Fourier expansion of:  $N(E) \backslash N(\mathbb{A}_E) \cong E \backslash \mathbb{A}_E \rightarrow \mathbb{C}$ ,  $n \mapsto f(ng)$ .

Get, when  $f \Rightarrow$  cuspidal, that

$$\forall g \in GL_2(\mathbb{A}_E), f(g) = \sum_{\gamma \in E^\times} W_{f, \psi} \left( \begin{pmatrix} \gamma & 0 \\ 0 & 1 \end{pmatrix} g \right) \quad (*)$$

with  $\psi: E \backslash \mathbb{A}_E \rightarrow \mathbb{C}^\times$  non-trivial character  
(used to define the FT)

and

$$W_{f, \psi}(g) = \int_{N(E) \backslash N(\mathbb{A}_E)} f(ng) \psi^{-1}(n) dn \quad \text{"first Fourier coefficient of } f \text{"}$$

$\in \text{Fun}(GL_2(\mathbb{A}_E), \mathbb{C})^{(N(\mathbb{A}_E), \psi)}$  "space of Whittaker functions"

The L-function of  $f$  can then be described as the Mellin transform of  $W_{f,\psi}$ . Useful perspective to generalize construction of L-functions to automorphic reps of  $GL_2, GL_n$  (Godement, Jacquet-Langlands).

Note that for any  $W$  Whittaker function, the formula (\*) with  $W_{f,\psi}$  replaced by  $W$  gives a cuspidal function  $f_W : \backslash GL_2(\mathbb{A}_E) \rightarrow \mathbb{C}$ .

with  $M = \begin{pmatrix} * & * \\ 0 & 1 \end{pmatrix} \subseteq GL_2$  mirabolic subgroup.

→ Can also be useful to construct automorphic forms.

Example (Shintani) Let  $\gamma = (\gamma_v)$  collection of conjugacy classes in  $GL_n(\mathbb{C})$  indexed by the places of  $E$ . To this one can attach a Whittaker function  $W_\gamma$  s.t.  $f_{W_\gamma} \in \text{Funcusp}(M(E) \backslash GL_2(\mathbb{A}_E), \mathbb{C})$  is right-invariant by maximal compact and Hecke eigenform with Hecke eigenvalue at  $v$  determined explicitly by  $\gamma_v$ . Moreover, unique with these properties.

Using this, the unramified part of the "Galois  $\rightarrow$  automorphic" direction of the Langlands conjecture can be rephrased as:

If there exists  $\sigma: \text{Gal } E \rightarrow \text{GL}_n(\overline{\mathbb{Q}_\ell})$  irreducible, everywhere unramified s.t.  $\forall$  place  $v$ ,  $\sigma_v = \sigma(\text{Frob}_v)$ , then the function  $f_\sigma := f_{W_\sigma}$  is left-invariant by  $\text{GL}_2(E)$  (hence a cuspidal Hecke eigenform for  $\sigma$ ).

One can prove this conjecture by geometric means, when  $E =$  function field (Drinfeld form)

$X$  smooth projective geom. curve over  $k = \mathbb{F}_q$  s.t.  $E = k(X)$ . Let:

- For  $i \geq 0$ ,  $\mathcal{E}_i =$  moduli stack of flat coherent sheaves of generic rank  $i$ .

- For  $\mathcal{E} =$  universal coherent sheaf on  $\mathcal{E}_1 \times X$ ,

let  $\mathcal{V} = \underline{\text{Hom}}(\mathcal{O}_X, \mathcal{E}) \supseteq \mathcal{V}^\circ =$  substack where the section is injective.  
 vector bundle on (an open substack of)  $\mathcal{E}_1$

• let  $\mathcal{V}^\vee = \underline{\text{Ext}}^1(\mathcal{E}, \omega_X)$

$\text{Bun}'_2^U =$  locus where the extension is a rank 2 v.b. (i.e., unprim) = moduli of pairs  $(\mathcal{F}, s)$ ,  $\mathcal{F}$  rank 2 v.b.,  $s: \omega_X \hookrightarrow \mathcal{F}$ .

Note that, as notation suggests,  $\mathcal{V}$  and  $\mathcal{V}^\vee$  are dual vector bundles over  $\mathcal{E}_1$ . The duality pairing is given by

$$\mathcal{V} \times \mathcal{V}^\vee \rightarrow \underline{\text{Ext}}^1(\mathcal{O}, \omega_X) \cong \mathbb{A}^1$$

The character  $\psi$  gives Artin-Schreier sheaf  $\mathcal{L}_\psi$  on  $\mathbb{A}^1$  and a Fourier equivalence: (Deligne, Laumon)

$$\text{Det}(\mathcal{V}, \overline{\mathbb{Q}}_\ell) \cong \text{Det}(\mathcal{V}^\vee, \overline{\mathbb{Q}}_\ell).$$

(Here,  $\psi: k \rightarrow \overline{\mathbb{Q}}_\ell^\times$ , induces  $\psi: E \setminus \mathbb{A}^1_E \rightarrow \overline{\mathbb{Q}}_\ell^\times$  by  $(a, v) \mapsto \psi(\sum_v t_{v/k} (\text{Res}(a_X \omega)))$   
 $\omega \in \mathcal{O}'_{E/k}$  of

Now,  $\sigma$  can be seen as a line sheaf on  $X$  (since  $\pi_1(X) = \text{Gal}_E^{\text{unr}}$ ).

Even better, can be promoted to  $\mathcal{L}_\sigma \in \text{Det}(\mathcal{E}_0, \overline{\mathbb{Q}}_\ell)$ .

Have a map  $\mathcal{V}^\circ \rightarrow \mathcal{E}_0$  (taking  $\psi$ -kernel of the rational).

Can pullback  $\mathcal{L}_\sigma$  along it and !-extend to  $\mathcal{V}$ .

Then Fourier transform and restrict to  $\text{Bun}'_2$ .

Get  $\text{Aut}'_{\sigma} \in \text{Dit}(\text{Bun}'_2, \overline{\mathbb{Q}})$ .

By taking trace of Frobenius, geometrizes the function  $f_{\sigma}$  mentioned above. Proving the conjecture amounts to showing that  $\text{Aut}'_{\sigma}$  descends to a "Hecke eigenstuff" on  $\text{Bun}_2$  along the forgetful map  $\text{Bun}'_2 \rightarrow \text{Bun}_2, (\mathbb{F}, s) \mapsto \mathbb{F}$ .

## (2) A local geometric analogue: $\ell$ -adic coefficients

In recent years, spectacular advances in using such geometric methods in the local setting.  
(Faltings-Scholze, Zhu)

From now on,  $E =$  local non-arch. field.

Attached to  $E$ , there is for each perfectoid space  $S$  over its residue field, a "family of curves"

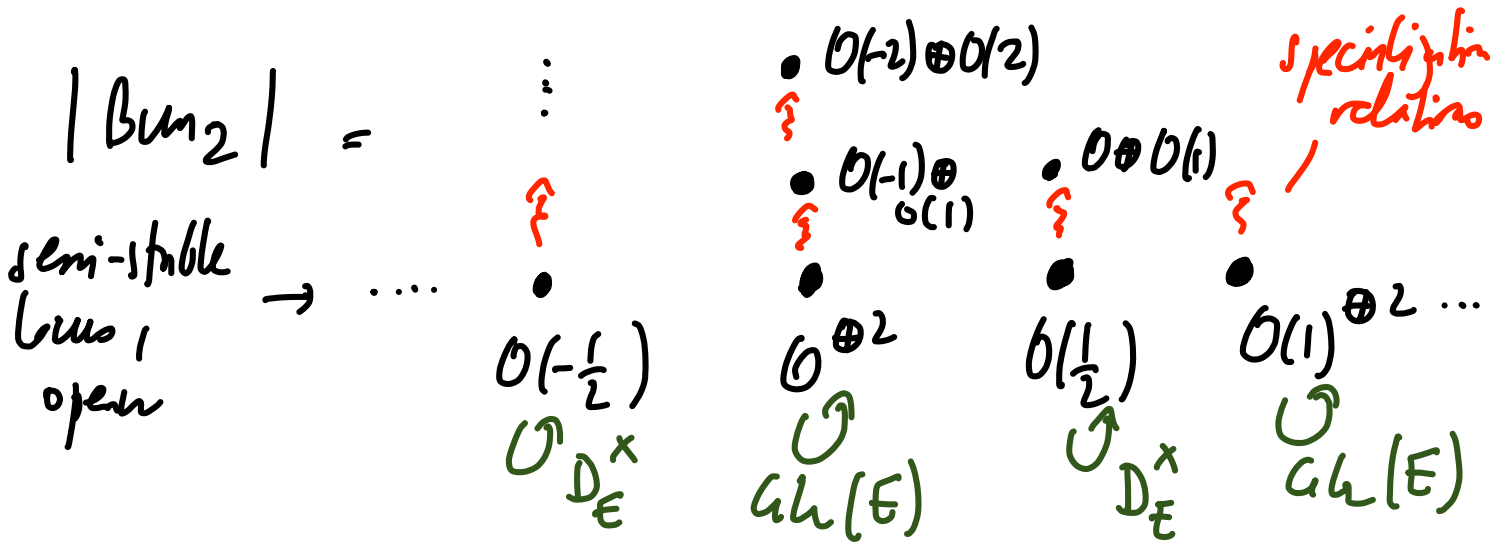
$X_{S, E}$  ("relative Faltings-Fantarc curve")

Can thus define:

•  $\text{Div}^1 : S \mapsto \left\{ \begin{array}{l} \text{effective relative Cartier} \\ \text{divisors of degree 1 on } X_{S, E} \end{array} \right\}$

$\ell$ -adic rep of  $W_E \hookrightarrow$  live  $\ell$ -adic sheaves on  $\text{Div}^1$ .

•  $Bun_2 : S \rightarrow$  groupoid of rank 2 vector bundles on  $X_{S,E}$ .



$\leadsto D(\text{smooth } \overline{\mathbb{Q}}\text{-rep of } GL(E)) \subseteq D_{\text{ét}}(Bun_2, \overline{\mathbb{Q}})$   
 $D(\text{smooth } \overline{\mathbb{Q}}\text{-rep of } D_E^x)$  fully faithful

Question: What does the above construction of Drinfeld(-Laurson) give in this context?

Key point: understand what the correct Fourier transform is in this setting. Earlier, had

$\underline{\text{Ext}}^1(\mathcal{O}_X, \omega_X) \cong A^1$  with its Artin-Schreier sheaf attached to  $\psi$ .  
 Serre duality

$\leadsto$  Fourier transform for sheaves of sections of flat coherent sheaves on  $X$ .

It. work with Anshuik: When  $X$  is the  
Fargues-Fontaine curve, replace  $\omega_X$  by  $\mathcal{O}_X$ . The  
stack of extensions of  $\mathcal{O}_X$  by  $\mathcal{O}_X$  is  $[* | \underline{E}]$ .  
Smooth character  $\psi: E \rightarrow \overline{\mathbb{Q}}^\times$  tautologically gives  
rise to rank 1 line sheaf on  $[* | \underline{E}]$ .

$\rightsquigarrow$  Fourier transform for Banach-Colmez  
spaces (= sheaves of sections of "flat" coherent  
sheaves on the FF curve).

Using this, can rerun the above construction  
in this new geometric setting.

We therefore fix  $\sigma: W_E \rightarrow \text{GL}_2(\overline{\mathbb{Q}})$  continuous  
irreducible Weil representation. Hope to produce  
 $\text{Ant}'_\sigma \in D(\text{Ban}'_2, \overline{\mathbb{Q}})$  descending to  $\text{Ant}_\sigma \in D(\text{Ban}_2, \overline{\mathbb{Q}})$   
whose restriction to the semi-stable locus relates to  
 $\sigma$  via LLC + JLC.

To see that this is not unreasonable, can  
specialize to the fiber of  $\text{Ban}'_2 \rightarrow \text{Ban}_2$   
over the locus where the v.b.F is isomorphic to  
 $\mathbb{O}^{\oplus 2}$  or  $\mathbb{O}(\frac{1}{2})$ :



• Case  $\mathcal{F} \approx \mathcal{O}^{\oplus 2}$ . This fiber is  $[*/M(E)]$   
 and the map the natural map  $[*/M(E)] \rightarrow [*/GL_2(E)]$

Diagram:

$$\begin{array}{ccc}
 [*/M(E)] \cong [*/M(E)] & & [E/E^*] \cong E^*/E^* \cong * \\
 \downarrow & \swarrow & \downarrow \\
 [*/GL_2(E)] & & [* / E^*]
 \end{array} \quad (1)$$

→ the sheaf on  $[*/M(E)]$  produced by the  
 Drinfeld construction is  $\text{C-ind}_{M(E)}^1 \psi$   
 = Kirillov model of any supercuspidal rep!

• Case  $\mathcal{F} \approx \mathcal{O}(\frac{1}{2})$ .

Diagram:

$$\begin{array}{ccc}
 \left[ \frac{BC(\mathcal{O}(\frac{1}{2})) \text{Isot}}{D_E^x} \right] \subseteq \left[ \frac{BC(\mathcal{O}(-1))}{E^*} \right] & \rightarrow & [* / E^*] \\
 \downarrow & \swarrow & \downarrow \\
 [* / D_E^x] & & \left[ \frac{BC(\mathcal{O}(1))}{E^*} \right] \cong \text{Div}^1
 \end{array} \quad (2)$$

Hence, geometric construction from  $\sigma$  of a  $D_E^x$ -equivariant sheaf on  $BC(\mathcal{O}(\frac{1}{2})) \text{Isot}$  which should be constant as a sheaf, given by the smooth rep of  $D_E^x$  attached to  $\sigma$  by LLC + JLC.

→ In particular, should have finite rank and can at least check it has the correct rank (Deligne, Carayol) using the adic GOS formula (Kamuro, Huber).

Rk: Construction of  $L_\sigma$  still has to be done, and needs genuinely new ideas to prove descent to  $\text{Bun}_2$ .

③ A local geometric analogue: towards mod  $p$  coefficients? (in progress)

Assume  $E \geq \mathbb{Q}_p$ .

The recent PhD thesis of Mann allows to define a category  $\mathcal{D}(\text{Bun}_2)$  of "mod  $p$  sheaves on  $\text{Bun}_2$ ". This category has a behaviour which is in between étale and coherent sheaves.

mod  $p$   $(\varphi, P)$ -modules  $\leftrightarrow$  dualizable objects (in dg  $\infty$ ) in  $\mathcal{D}(\text{Div}')$ .  
(Lubin-Tate version)

Is there an analogue of Drinfeld construction in this setting too? (Question independently asked and studied by Fargues)

Fourier theory has to work differently. Geometric objects stay the same (Banach-Colony spaces), but coefficients have changed, and it now seems the kernel of the FT should be the rank 1 object in  $D([* / BC(O(d))])$ ,  $d = [E : \mathbb{Q}_p]$ , induced by  $BC(O(d)) \simeq \tilde{G}_m$ .

Hence, in Drinfeld construction, the diagram appearing when specializing to the fiber of  $Bun_2$  over the locus where the bundle is isomorphic to  $O(d)^{\oplus 2}$  now looks like a mix between

(1) and (2):

$$\begin{array}{ccc}
 [* / M(E)] \simeq [* / M(E)] & \xrightarrow{[BC(O(d)) / E^*]} & Div^d \\
 \downarrow & \searrow & \\
 [* / GL(E)] & & [* / E^*]
 \end{array}$$

To  $\sigma: GL_E \rightarrow GL(\overline{\mathbb{F}}_p)$  irreducible

$\Leftrightarrow$  irreducible rank 2  $(\varphi, \Gamma)$ -module seen in  $Div^1$ , can thus attach a rep of the mirabolic subgroup. This is (should be...) Colony construction when  $d=1$ , i.e.  $E = \mathbb{Q}_p$ .

## ④ Whittaker models

Global/local Whittaker models: embedding of a generic (e.g. tempered) representation in the space of Whittaker functions (seen as a rep of the group by right-translation action). Dually: any such representation admits a surjection from the compact induction of  $\psi$  from upper unipotent to  $GL_n =:$  the Whittaker rep.

Geometric Langlands perspective on this: for any choice of base (global/local) and coefficients ( $\ell$ -adic, mod  $p$ , ...), should have an equivalence

$$D(\text{Bun}_n) \simeq \text{Ind Coh}_{\text{MfP}} \left( \begin{array}{l} \text{moduli of} \\ L\text{-parameters} \\ \text{for } GL_n \end{array} \right)$$

compatible with Hecke actions and many other structures. We do not try to make the statement more precise here.

Let us point out that if we restrict to the "tempered part" on the LHS (in particular, disregard anything

non generic), can just pretend that one considers  $Q_{\text{glob}}$  on the RHS.

Let  $\mathcal{W} = \underline{\text{Whittaker sheaf}} = \text{image of structure sheaf on the stack of } L\text{-parameters.}$

Since the structure sheaf maps surjectively on "the" skyscraper sheaf at any point of the stack of  $L$ -parameters,  $\mathcal{W}$  functions as a geometric incarnation of the Whittaker representation.

In fact,  $\exists$  general guess for what  $\mathcal{W}$  should be: Here for  $GL_2$  for simplicity. Consider:

$$\mathcal{N} = \begin{array}{c} \text{stack of extensions} \\ \text{of } \mathbb{Q}_X \text{ by } \omega_X \end{array} \xrightarrow{\quad} \text{Bun}_2$$

$\downarrow f$

(send the extension to the underlying bundle)

Note that in any of the three cases discussed above,  $\mathcal{N}$  is where the kernel of the Fourier transform lives. Its image by  $f!$  should be  $\mathcal{W}$ .

To finish, let us see what this says in the local setting:

- local setting,  $\ell$ -adic coefficients: gives  $\mathcal{W} = i_{1,!} (c\text{-ind}_{N(\mathbb{Q}_p)}^{G_h(\mathbb{Q}_p)} \psi)$ . In particular, fiber at  $\text{Bun}_2^1 = [* / G_h(\mathbb{Q}_p)] \underset{\text{open}}{\subseteq} \text{Bun}_2$  is the Whittaker representation.

- local setting, mod  $p$  coefficients: in this context, there is as far as I know no theory of the Whittaker model (basic problem: no  $\psi$ !). This is "explained" by the shape of  $\mathcal{W}$ : it is supported on the stratum given by  $\mathcal{O} \oplus \mathcal{O}(1)$ , hence is zero in restriction to  $\text{Bun}_2^1$ .